

Precision physics with polarised beams

Roman Pöschl

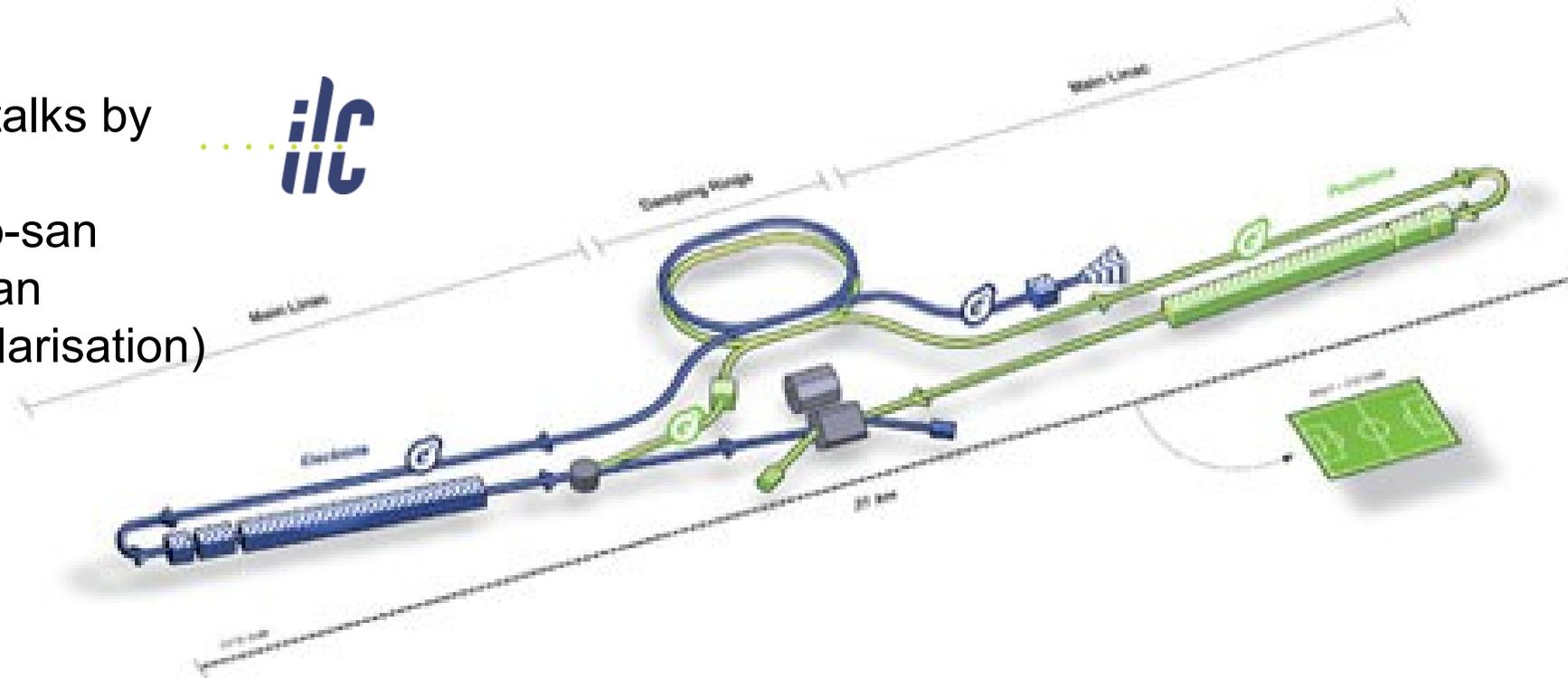


HKUST Jockey Club Hong Kong – January 2019

See also talks by



Michizono-san
Yokoya-san
(Beam polarisation)



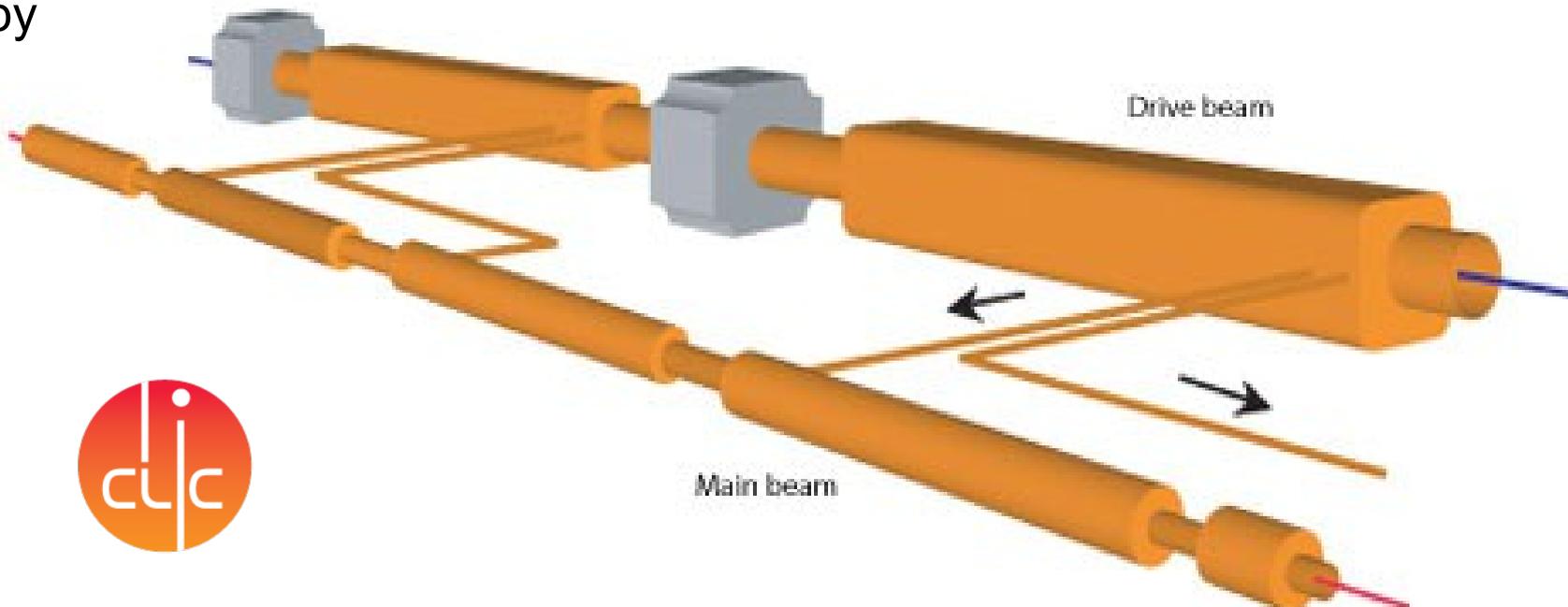
Energy: 0.1 - 1 TeV
Electron (and positron)
polarisation
TDR in 2013
+ DBD for detectors
Footprint 31 km

Initial Energy 250 GeV – Footprint ~20km

Strong effort by Japanese Community to host ILC – Political decisions expected ... daily

See also talk by

Andrea Latina



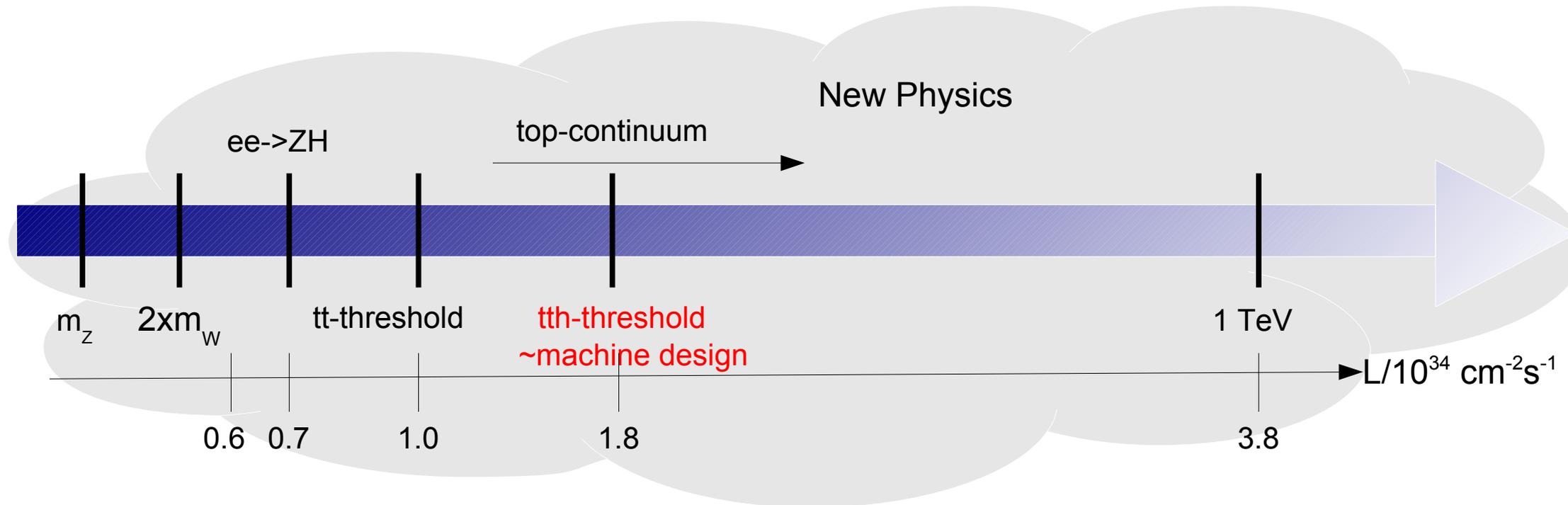
Energy: 0.4 - 3 TeV

CDR in 2012

Footprint 48km

Initial Energy 380 GeV

Possible future project at CERN



All Standard Model particles within reach of LC

- High precision tests of Standard Model over wide range to detect onset of New Physics

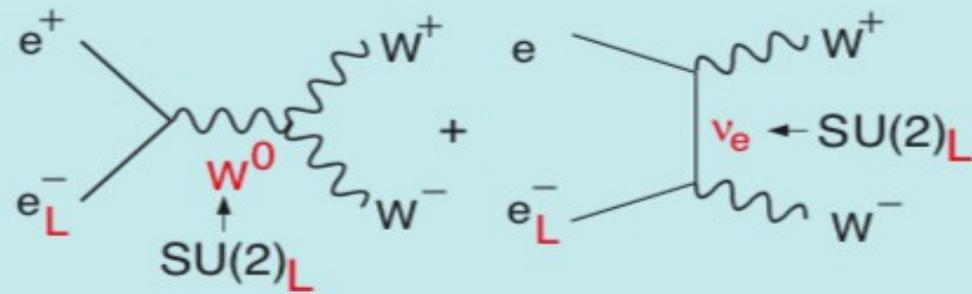
Machine settings can be “tailored” for specific processes

- Centre-of-Mass energy
- Beam polarisation

$$\sigma_{P,P'} = \frac{1}{4} [(1 - PP')(\sigma_{LR} + \sigma_{RL}) + (P - P')(\sigma_{RL} - \sigma_{LR})]$$

Background free searches for BSM through **beam polarisation**

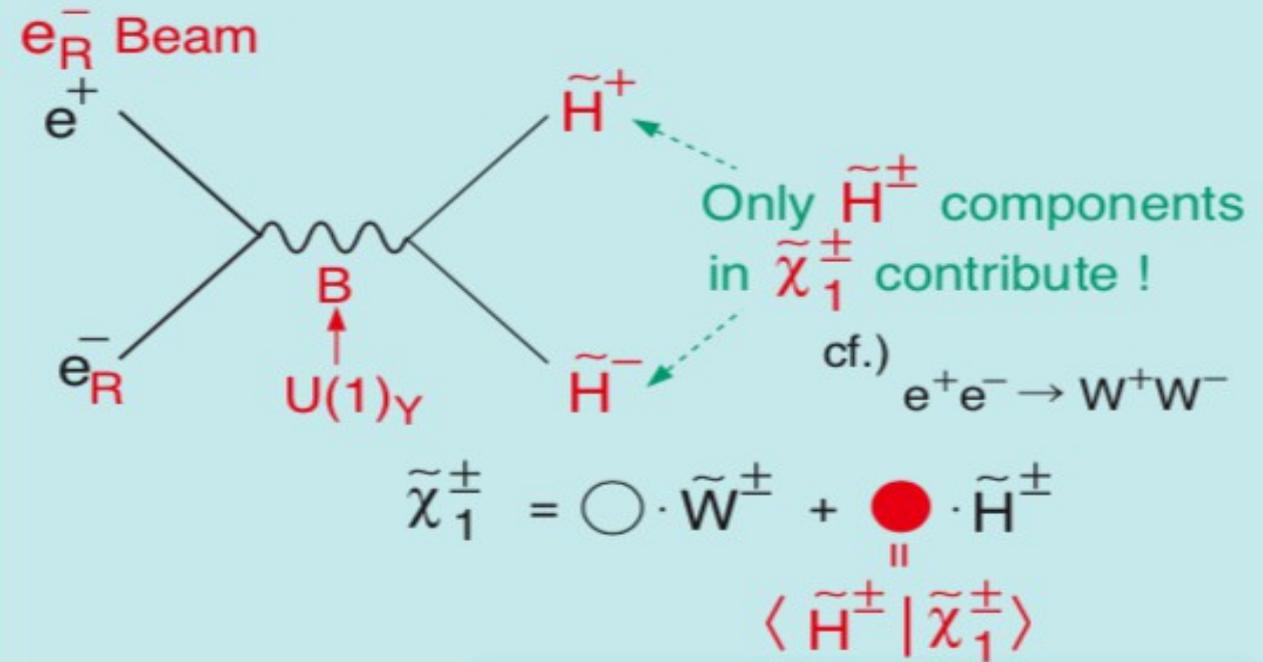
W^+W^- (Largest SM BG in SUSY searches)



In the symmetry limit, $\sigma_{WW} \rightarrow 0$ for e_R^- !

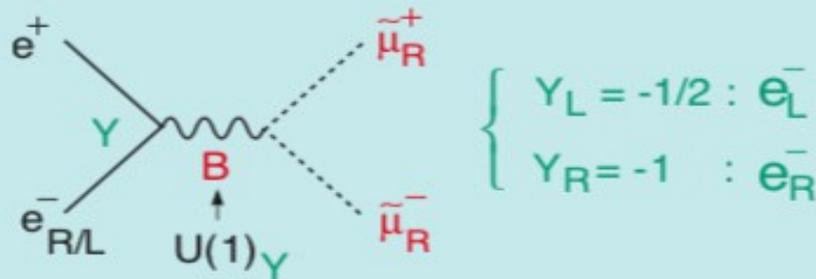
BG Suppression

Chargino Pair



Decomposition

Slepton Pair

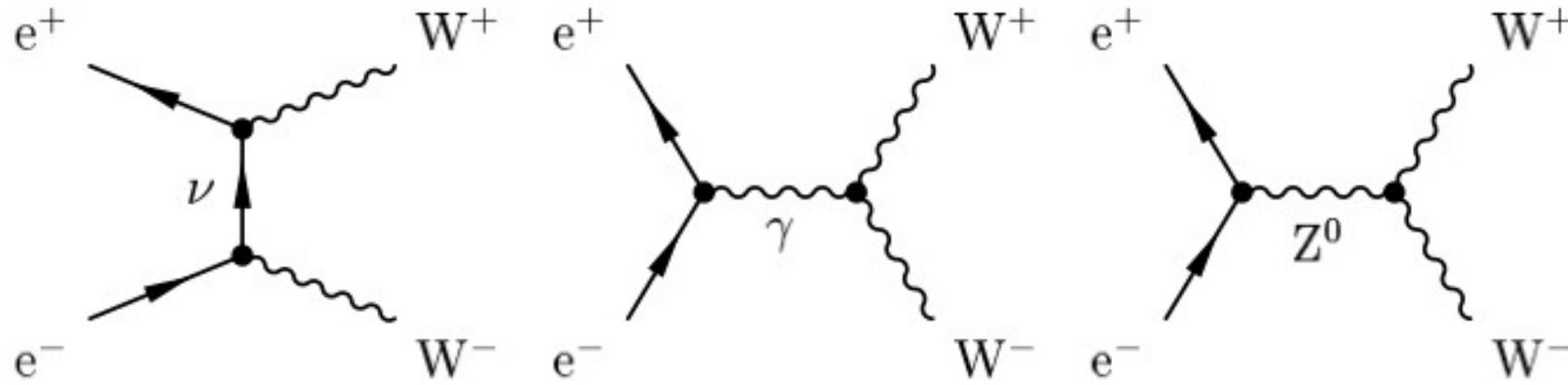


In the symmetry limit, $\sigma_R = 4 \sigma_L$!

WW-fusion Higgs Prod.



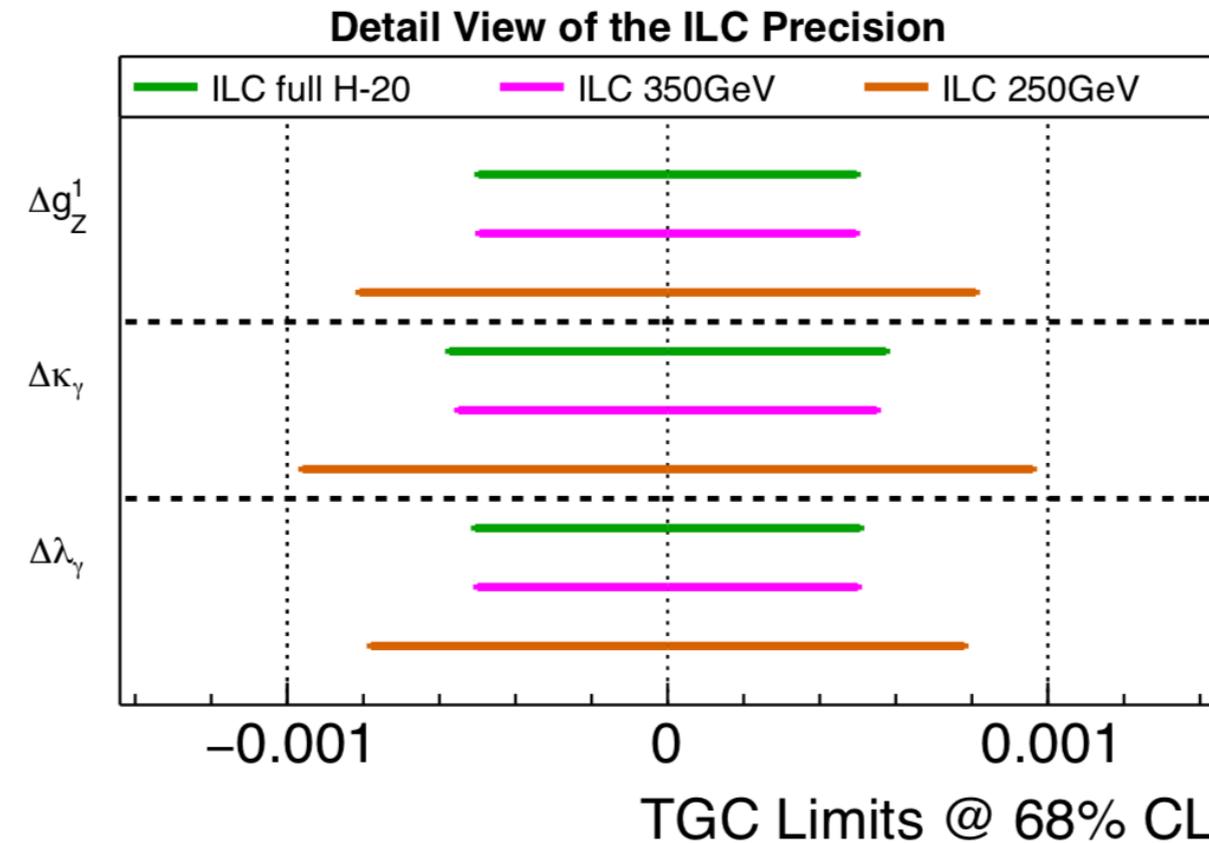
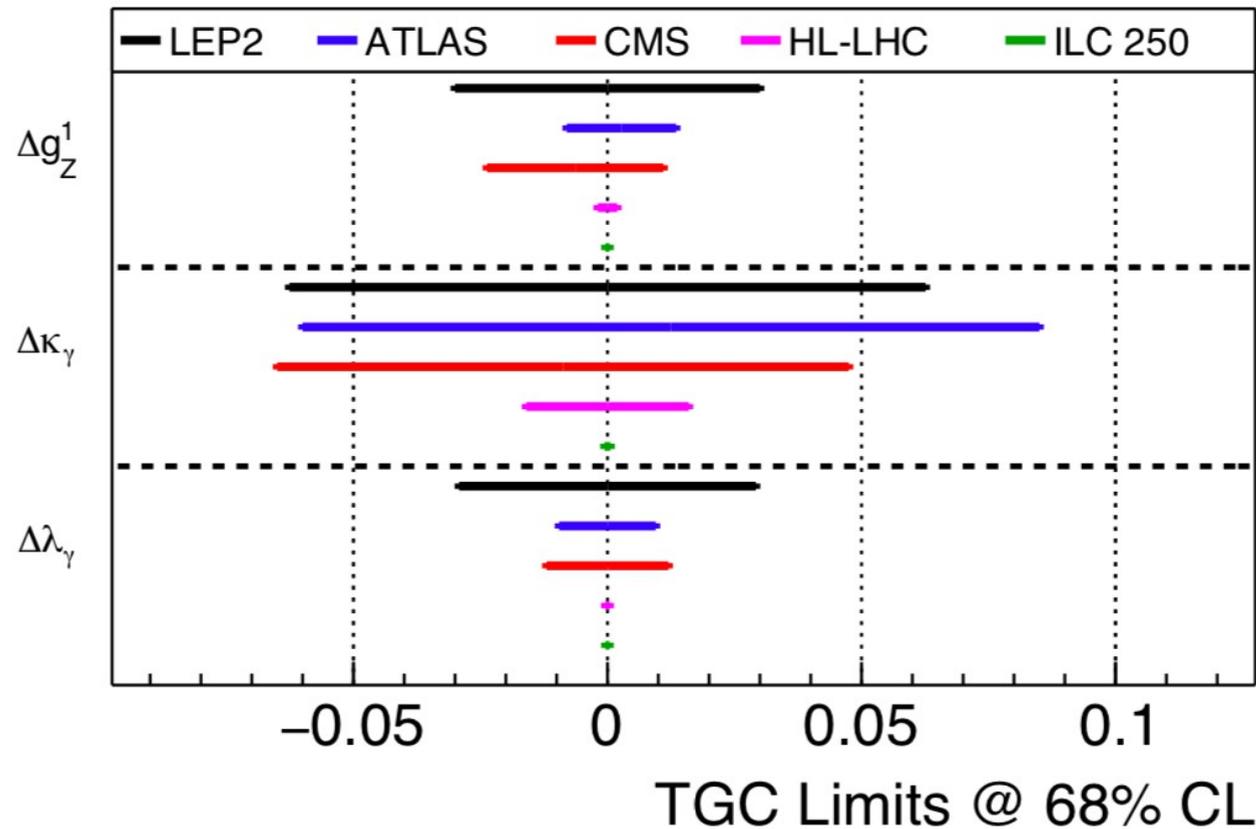
Signal Enhancement



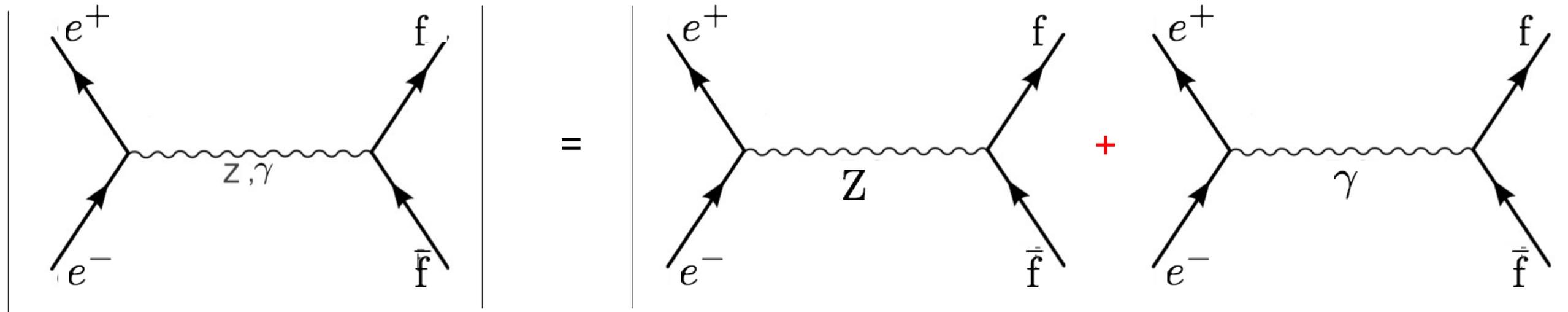
- Sensitivity to triple and quartic gauge Boson couplings (TGC and QGC)
- Observables depend strongly on beam polarisation

=> Enrich different helicity modes of W
=> Disentangling of couplings to Z and γ
=> in situ measurement of beam polarisation (and luminosity)

Limits on Triple Gauge Couplings @ 250 GeV



Cross section $e^+e^- \rightarrow f\bar{f}$



Interference between individual amplitudes of γ and Z exchange

$$\mathcal{M}_Z = -\frac{\sqrt{2}G_F M_Z^2}{s - M_Z^2} \left[\bar{f} \gamma^\rho (c_V^f - c_A^f \gamma^5) f \right] g_{\rho\sigma} \left[\bar{e} \gamma^\sigma (c_V^e - c_A^e \gamma^5) e \right]$$

$$\mathcal{M}_\gamma = -\frac{e^2}{s} (\bar{f} \gamma^\nu f) g_{\mu\nu} (\bar{e} \gamma^\nu e)$$

$$g_L^f = c_V^f + c_A^f$$

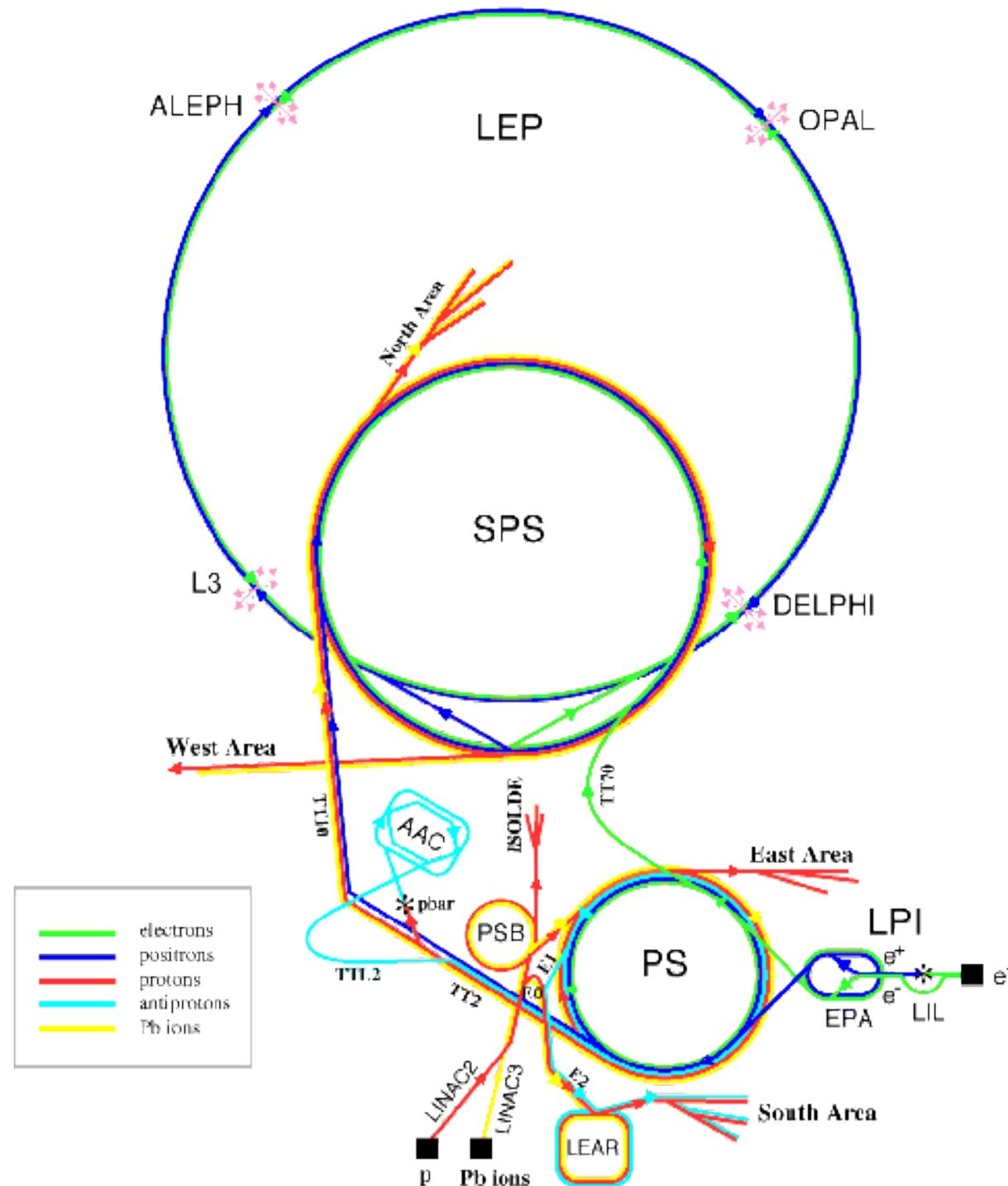
$$g_R^f = c_V^f - c_A^f$$

Differential cross section:

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} \left[A_0(1 + \cos^2\theta) + A_1 \cos\theta \right] \left\{ \begin{array}{ll} \sim (1 + \cos^2\theta) & \text{'Usual' Vector current, symmetric in } \cos\theta \\ \sim \cos\theta & \text{Axial Vector current, asymmetric in } \cos\theta \end{array} \right.$$

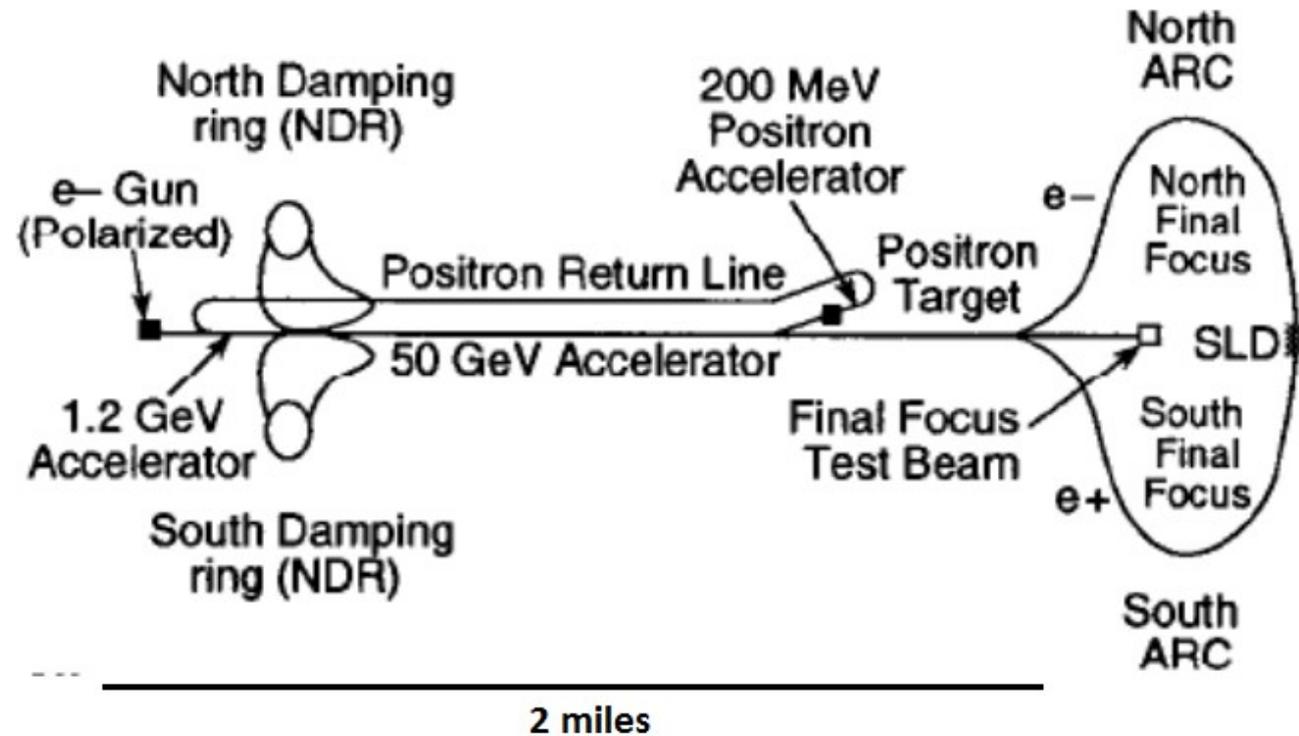
Weak interaction introduces forward backward asymmetry

=> Asymmetry is intrinsic to electroweak processes!!!



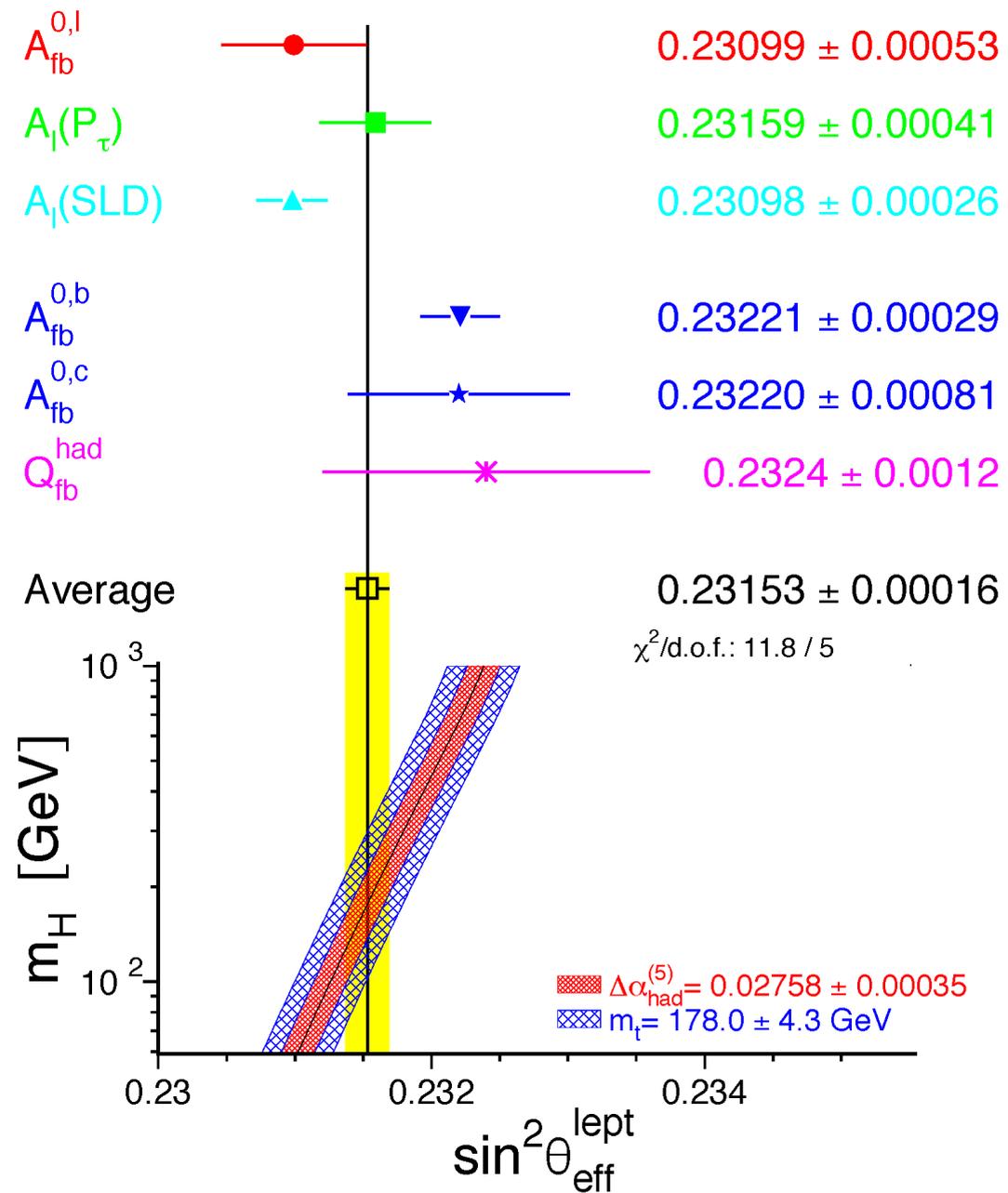
Large Electron Positron Collider – LEP:

- Circular electron positron collider
- Centre of mass energies m_Z – 209 GeV
- Operated at CERN between 1989 and 2000
- No beam polarisation but high luminosity at
 - four interaction points
 - Around 10M Z events collected
- No beam polarisation



SLAC Linear Collider – SLC:

- Linear electron positron collider
- Centre of mass energy m_Z
- Operated at SLAC between 1992 and 1998
- Electron beam polarisation 90%
- One single interaction point
 - Around 400k Z events collected



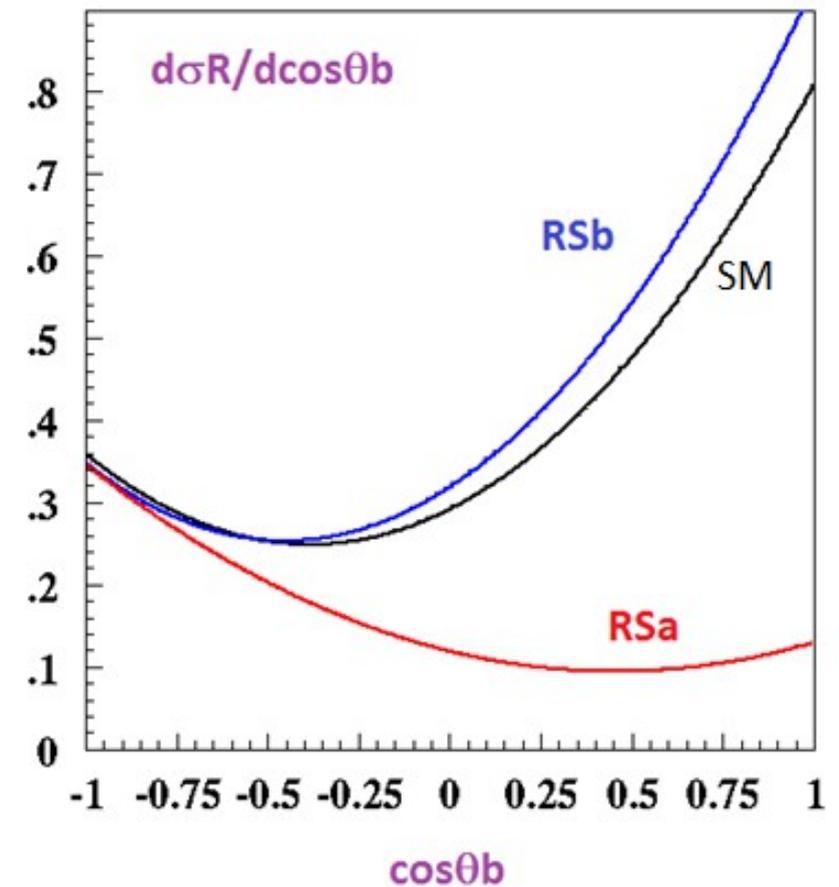
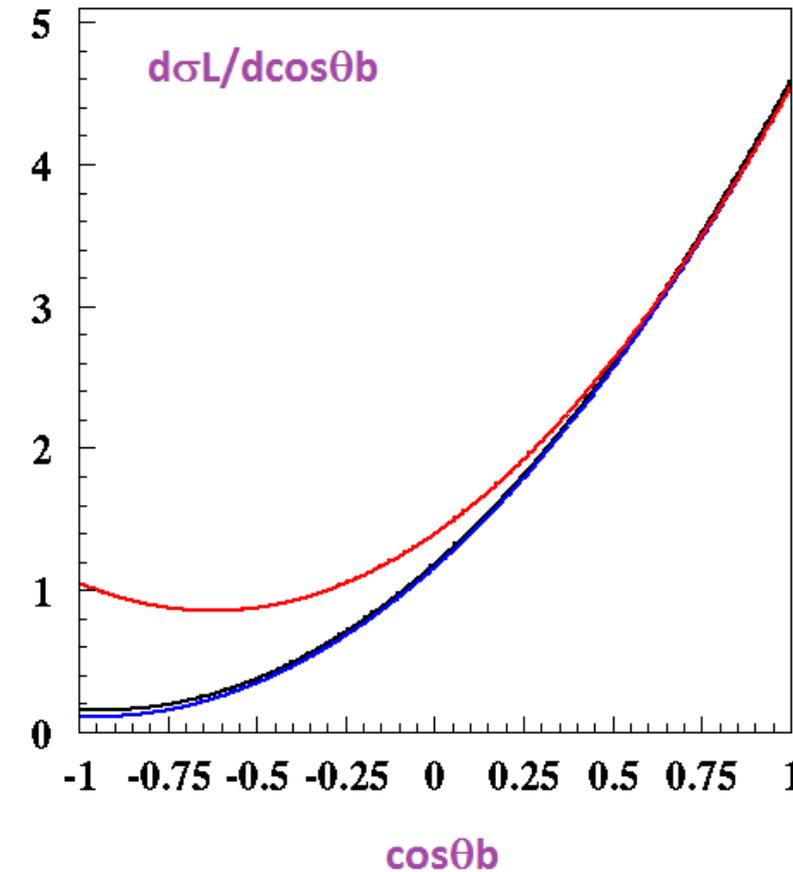
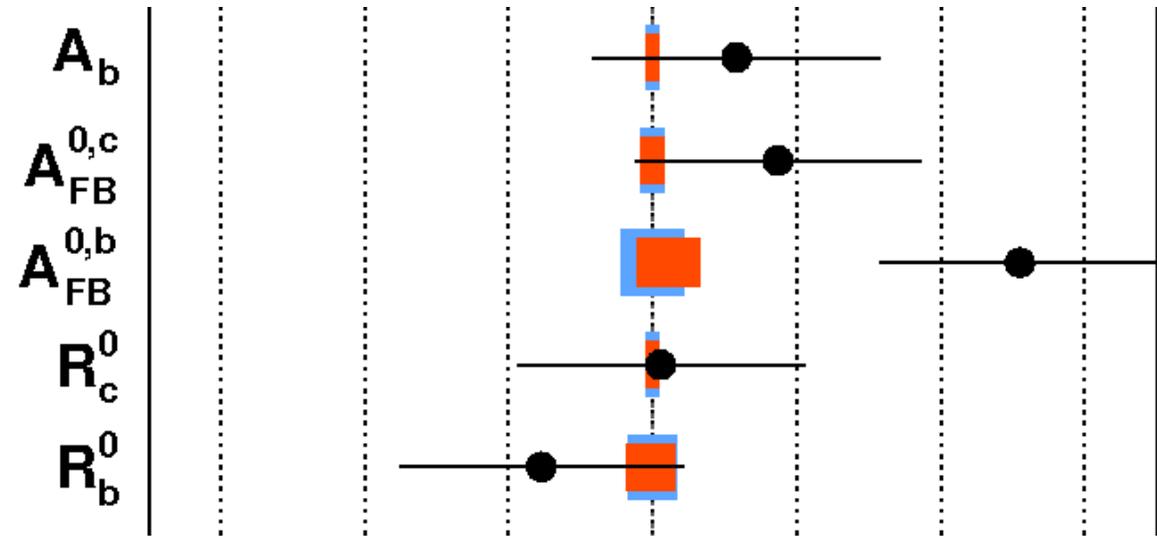
- Most precise single Individual determination of $\sin^2 \theta_{\text{eff}}^l$ from SLC
 - Left-right asymmetry of leptons
- Most precise measurement of $\sin^2 \theta_{\text{eff}}^l$ from forward backward asymmetry A_{FB}^b in $ee \rightarrow bb$ at LEP

Two lessons:

- Most precise determinations of $\sin^2 \theta_{\text{eff}}^l$ differ significantly
 - Cries for verification
 - Beam polarisation can match up for luminosity

$\sim 3\sigma$ in heavy quark observable A_{FB}^b

ee- \rightarrow bb@250 GeV



- Is tension due to underestimation of errors or due to new physics?

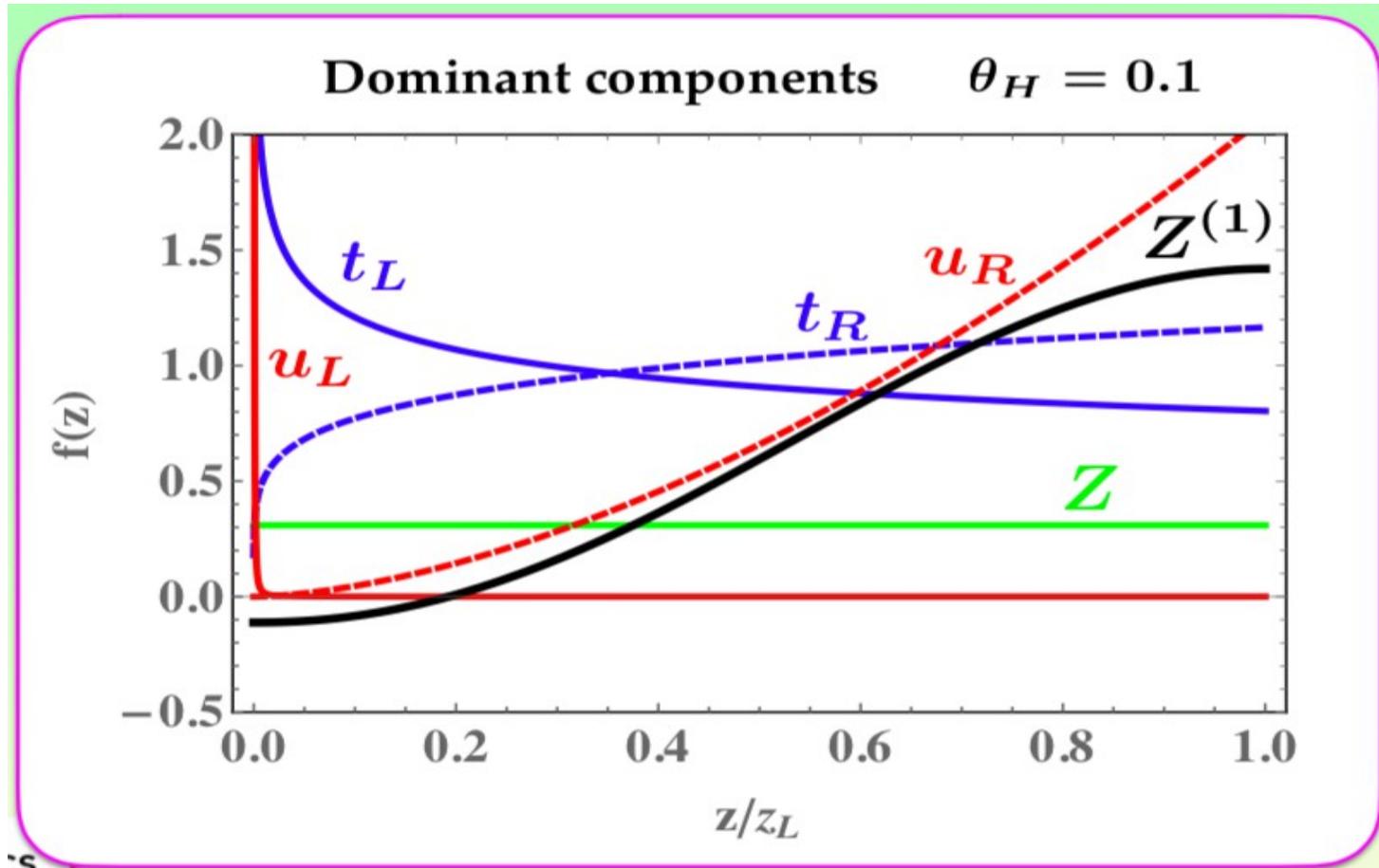
- High precision e+e- collider will give final word on anomaly

Randall Sundrum Models Djouadi/Richard '06

- In case it will persist polarised beams will allow for discrimination between effects on left and right handed couplings (Remember $Zb_l b_l$ is protected by cross section)

- Note that also B-Factories report on anomalies IAS 2019

Randall Sundrum Models imply arrangement of fermion wave functions in (warped) extra dimension
The more overlap on IR-Brane the larger the interaction

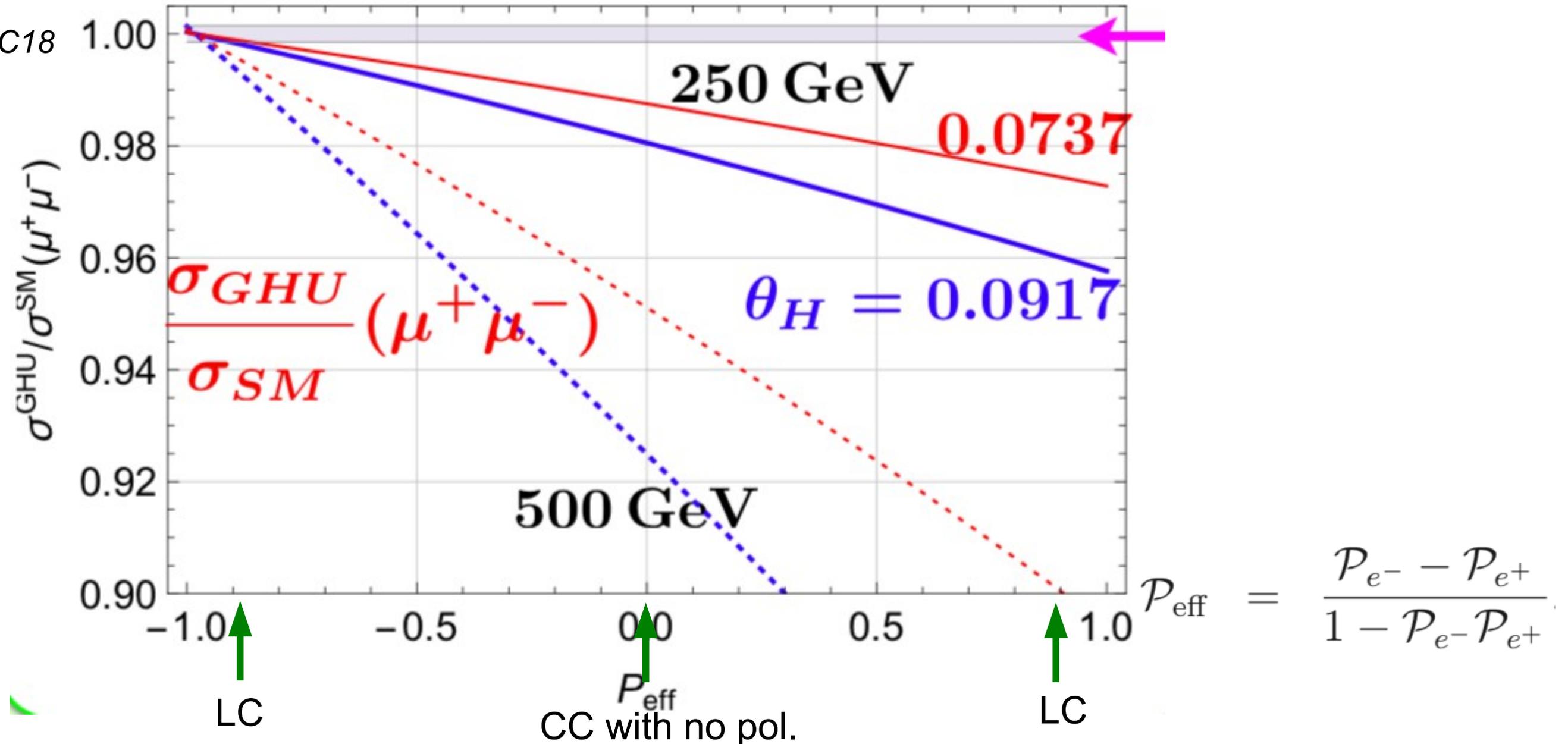


	SM: Z		$Z^{(1)}$		$Z_R^{(1)}$		$\gamma^{(1)}$	
	Left	Right	Left	Right	Left	Right	Left	Right
ν_e			-0.183	0	0	0	0	0
ν_μ	0.5	0	-0.183	0	0	0	0	0
ν_τ			-0.183	0	0	0	0	0
e			0.099	0.916	0	-1.261	0.155	-1.665
μ	-0.2688	0.2312	0.099	0.860	0	-1.193	0.155	-1.563
τ			0.099	0.814	0	-1.136	0.155	-1.479
u			-0.127	-0.600	0	0.828	-0.103	1.090
c	0.3458	-0.1541	-0.130	-0.555	0	0.773	-0.103	1.009
t			0.494	-0.372	0.985	0.549	0.404	0.678
d			0.155	0.300	0	-0.414	0.052	-0.545
s	-0.4229	0.0771	0.155	0.277	0	-0.387	0.052	-0.504
b			-0.610	0.186	0.984	-0.274	-0.202	-0.339

GHU Model:

- Interaction of right handed (light) fermions -> **Heavy and light fermion effect**
- Interaction of left and right handed heavy quarks
- Note also asymmetry in couplings to $\gamma^{(1)}$ => **F1Ay ≠ 0**

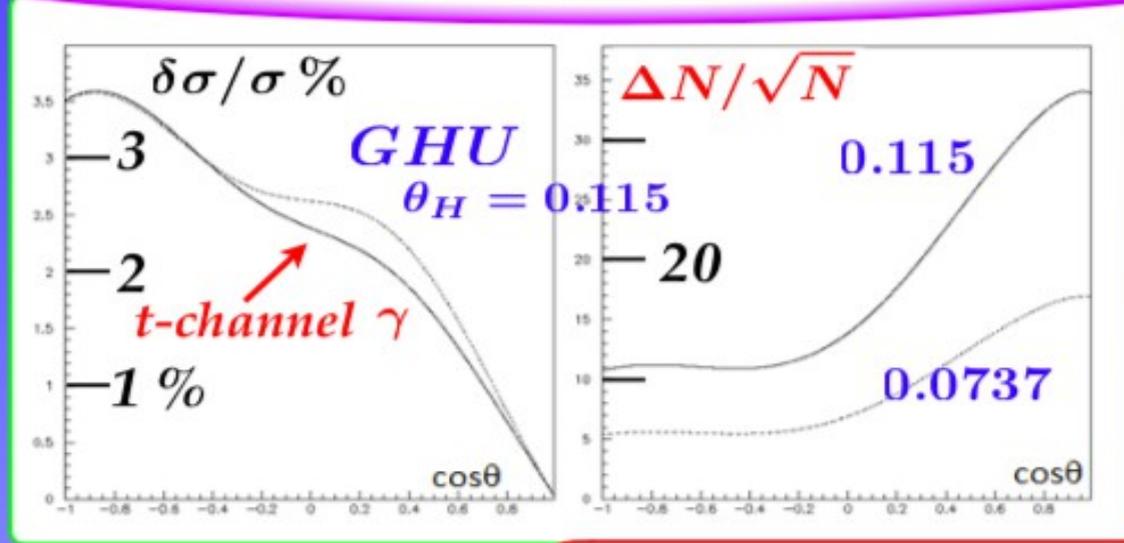
Y. Hosotani, Top@LC18



- Visible effects for $\mathcal{P}_{eff} \geq -0.5$
- LC would add two points that ideally are complemented by a point from Circular Collider
- Huge amplification of effect at higher centre-of-mass energies

top2018-hosotani.pdf (page 17 of 22)

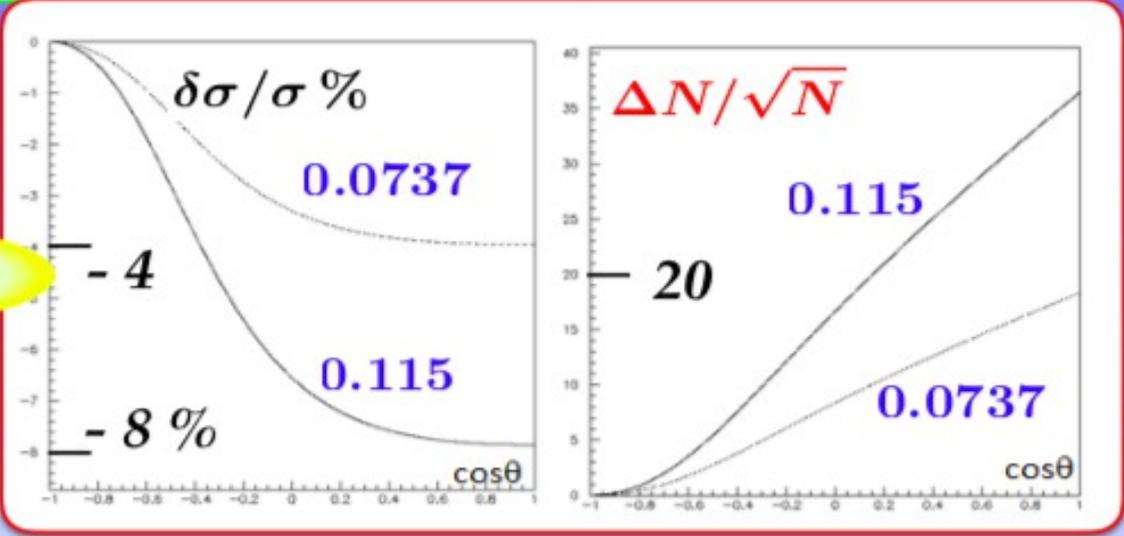
Bhabha scattering $e^+e^- \rightarrow e^+e^-$



F. Richard, 1804.02846

ILC 250
 2000 fb^{-1}
cos θ bin width 0.1

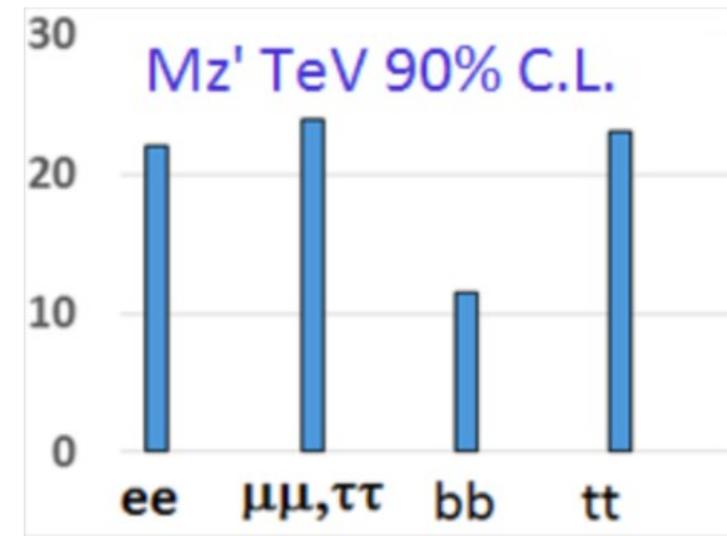
$e^+e^- \rightarrow \mu^+\mu^-$



GHU is example for model that implies

modifications of

“well known” couplings



Impressive mass reach already at ILC 250 GeV

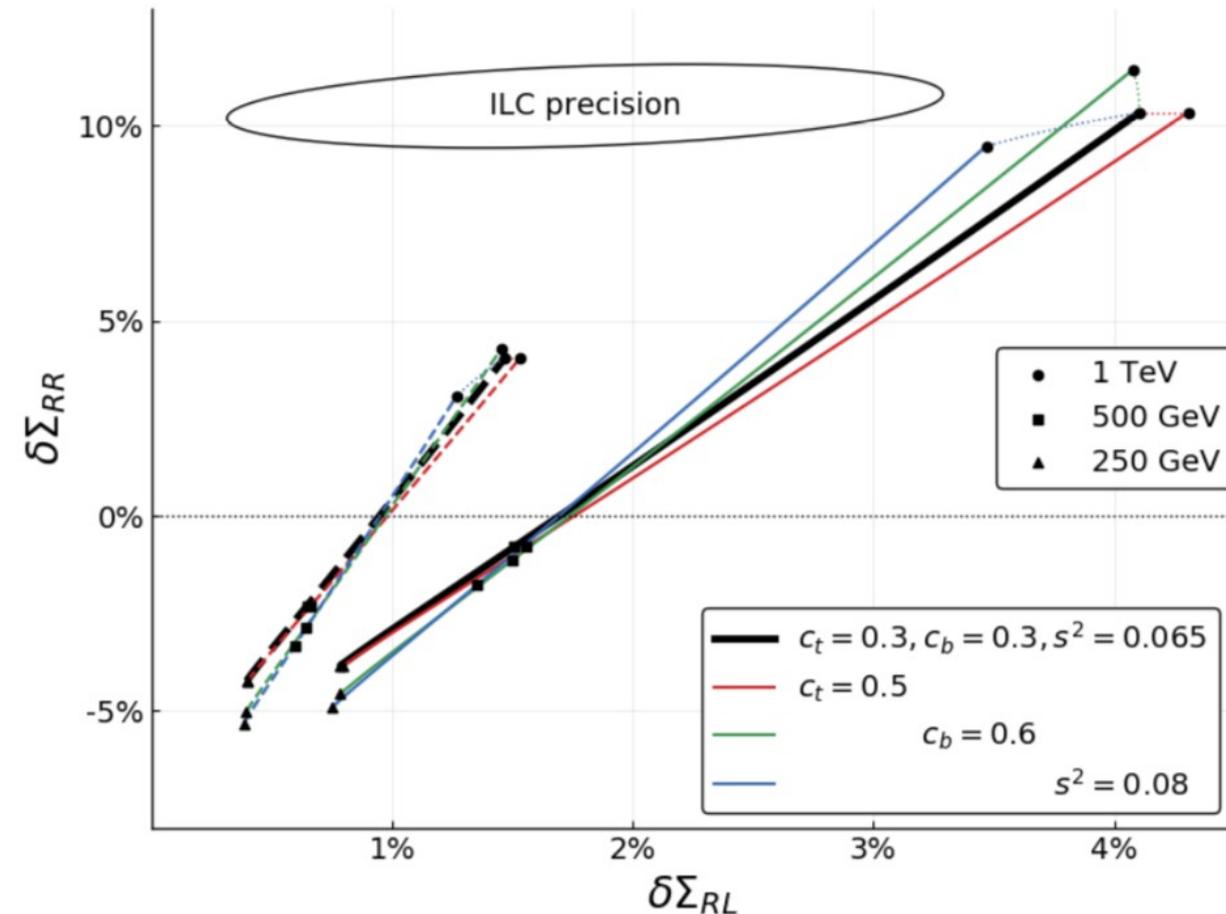
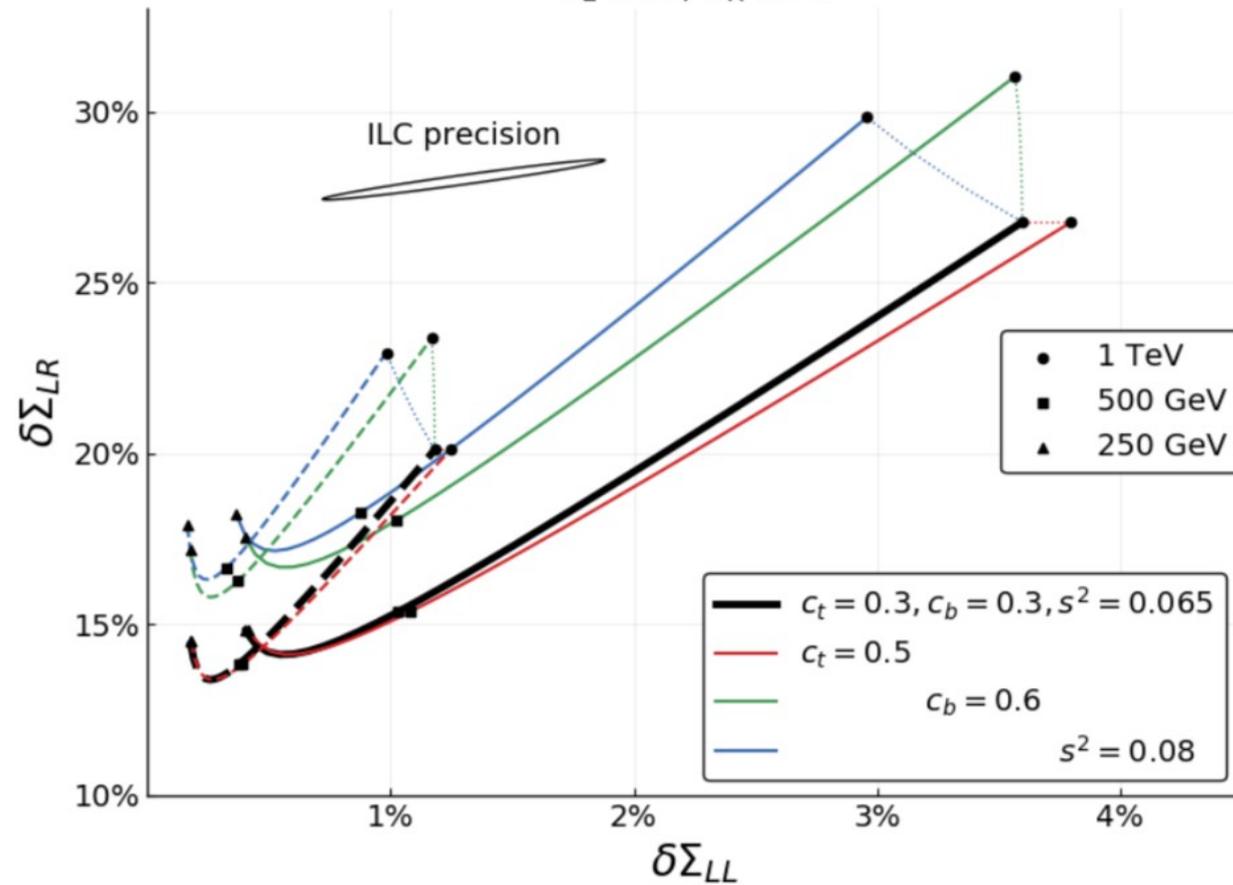
e.g.
$$\frac{m_b^2}{m_t^2} = \frac{1}{2} \tan^2 \theta_b \left(\frac{1 + 2c_b}{1 + 2c_t} \right) \left(\frac{z_0}{z_R} \right)^{2c_b - 2c_t}$$

In short, mixing is consequence of arrangement of **heavy quarks** in 5D multiplets

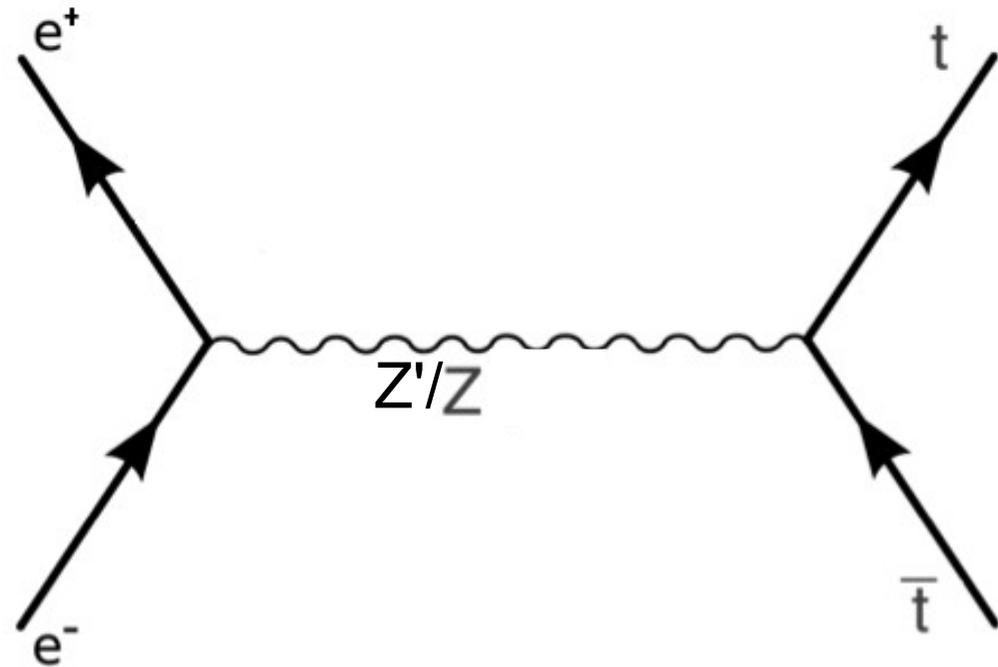
$$\frac{d\sigma}{d\cos\theta}(e_L^- e_R^+ \rightarrow b\bar{b}) = \Sigma_{LL}(s) (1 + \cos\theta)^2 + \Sigma_{LR}(s)(1 - \cos\theta)^2$$

$$\frac{d\sigma}{d\cos\theta}(e_R^- e_L^+ \rightarrow b\bar{b}) = \Sigma_{RL}(s) (1 - \cos\theta)^2 + \Sigma_{RR}(s)(1 + \cos\theta)^2$$

b_L in **5**, b_R in **4**

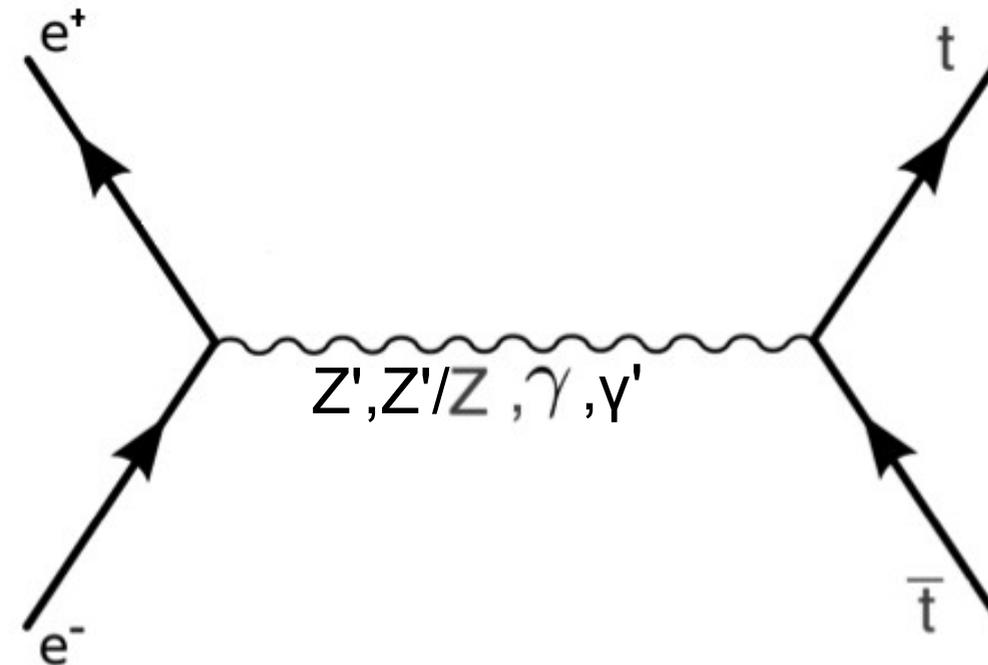


On the Z-pole



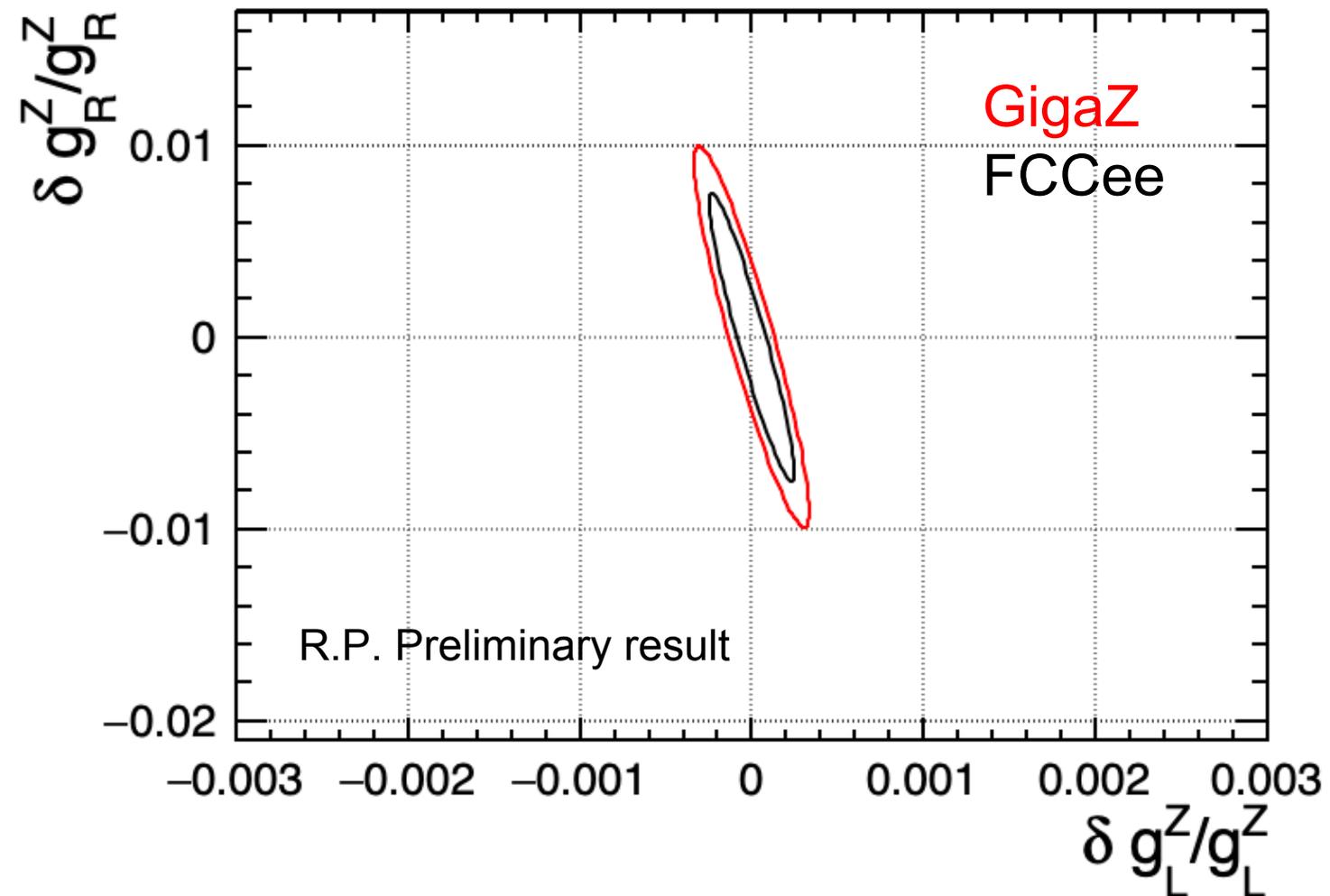
- Sensitivity to Z/Z' mixing
- Sensitivity to vector and tensor
- Couplings of the Z
 - (the photon does not “disturb”)

Above the Z-pole



- Sensitivity to interference effects of Z and photon!!
 - There is no reason to assume that the photon is standard model like, which is a model dependent assumption in EFT fits!!!

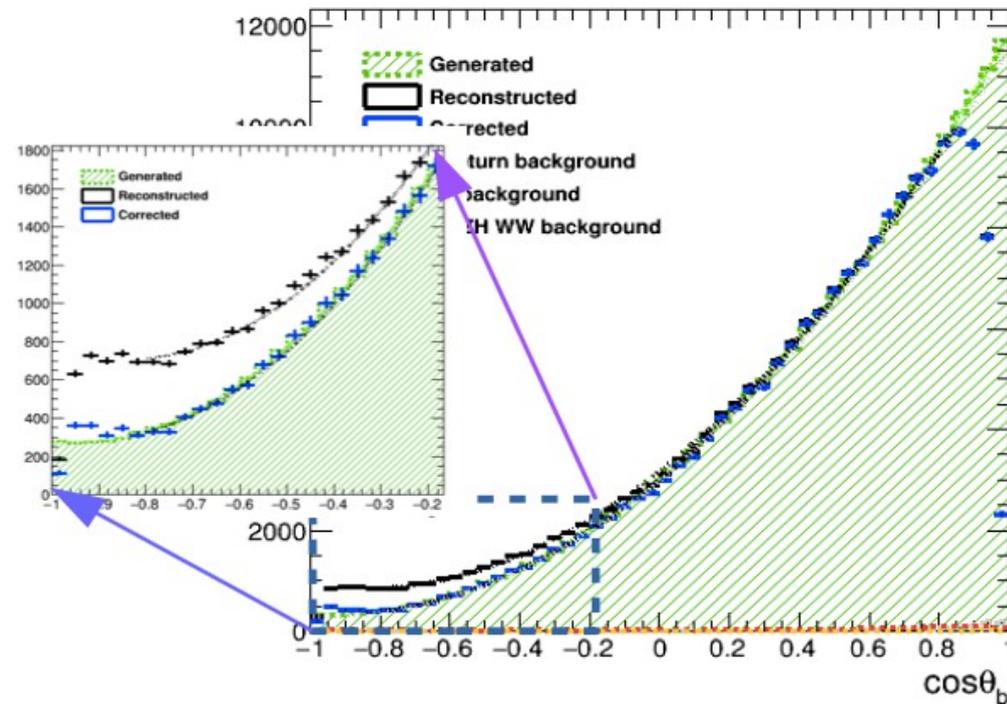
b-quark couplings on the Z pole



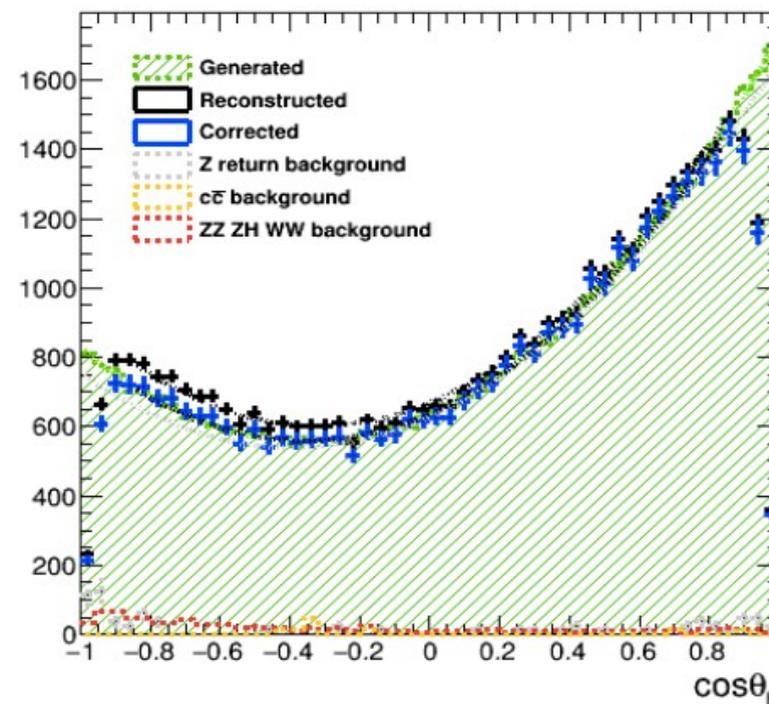
- **GigaZ results extrapolated from LEP1**
 - Taking into account excellent b-tagging at e^+e^- detectors but also split of lumi due to polarisation
 - Resulting statistical error multiplied by two to account for systematics
- **FCCee using optimistic estimation of systematic error**
 - See recent CDR
- **Upshot:**
 - Precise tests for Z/Z' mixing on Z pole
 - Dominant effect through statistics
 - Polarisation (100% e^- only) improves result by $\sim 30\%$
 - Better control of systematic errors (see backup)
 - GigaZ competitive

$$\sqrt{s} = 250 \text{ GeV} \quad L = 250 \text{ fb}^{-1}$$

$$e_L^- e_R^+ \rightarrow b\bar{b} \qquad e_R^- e_L^+ \rightarrow b\bar{b}$$



$$A_{fb}^{rec} / A_{fb}^{gen} = 100.7\% \pm 0.62\%$$



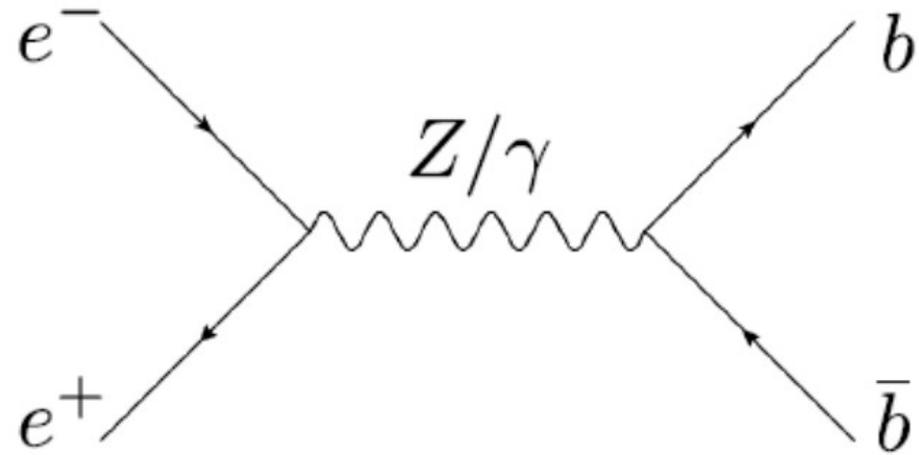
$$A_{fb}^{rec} / A_{fb}^{gen} = 104.9\% \pm 2.25\%$$

Full simulation study (with ILD concept), Benchmark reaction for 250 GeV running

- Experimental challenge: Measurement of b-quark charge on event-by-event basis

Long lever arm in $\cos \theta_b$ to extract form factors or couplings

$$\frac{d\sigma^I}{d\cos\theta} = S^I (1 + \cos^2 \theta) + A^I \cos \theta \qquad I = L, R \quad \begin{array}{l} \text{Form factors/couplings} \\ \text{from S and A} \end{array}$$



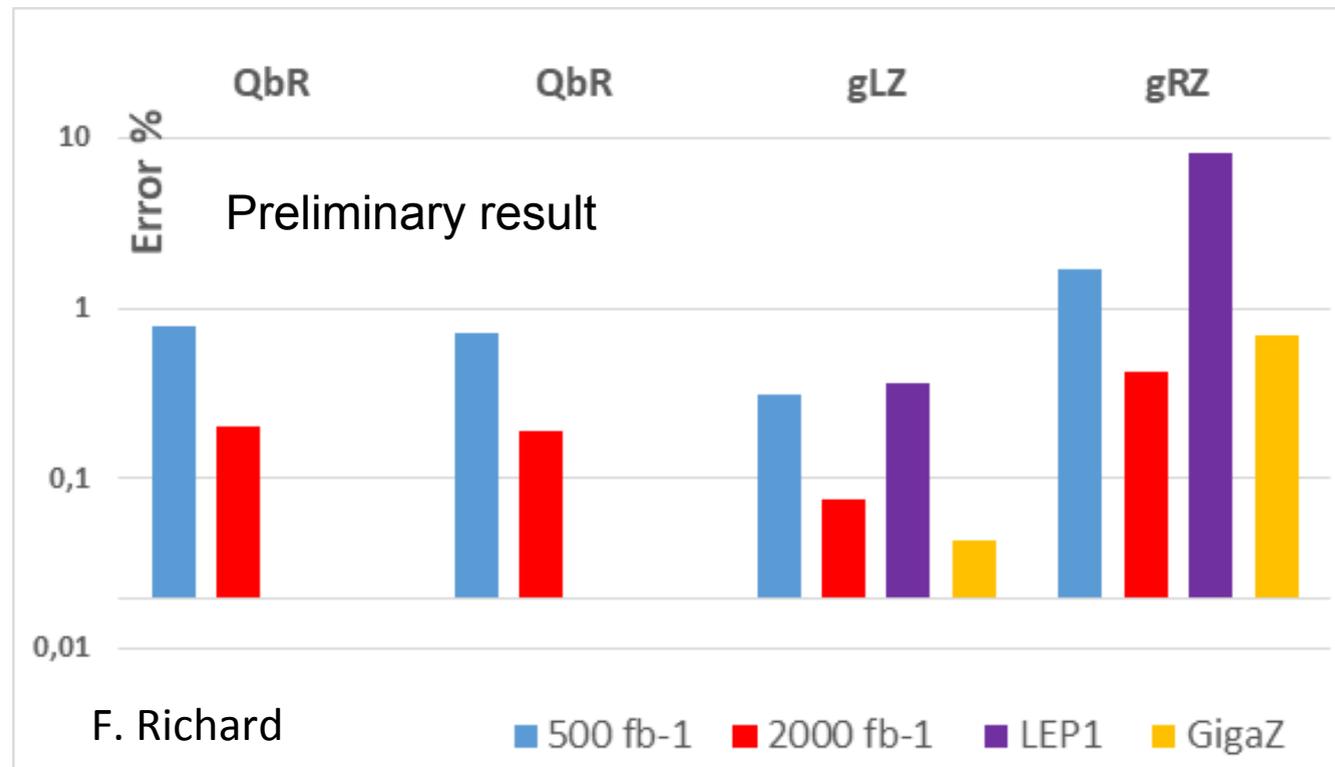
See also [1709.04289](#)

- ILC measurements with beam polarisation provide model independent access to photon and Z couplings (or vector and tensor couplings)
- $\delta g_{RZ}/g_{RZ} \sim 2\%$ sufficient to confirm at $>5\sigma$ or to discard the LEP1 effect which is at the 25% level

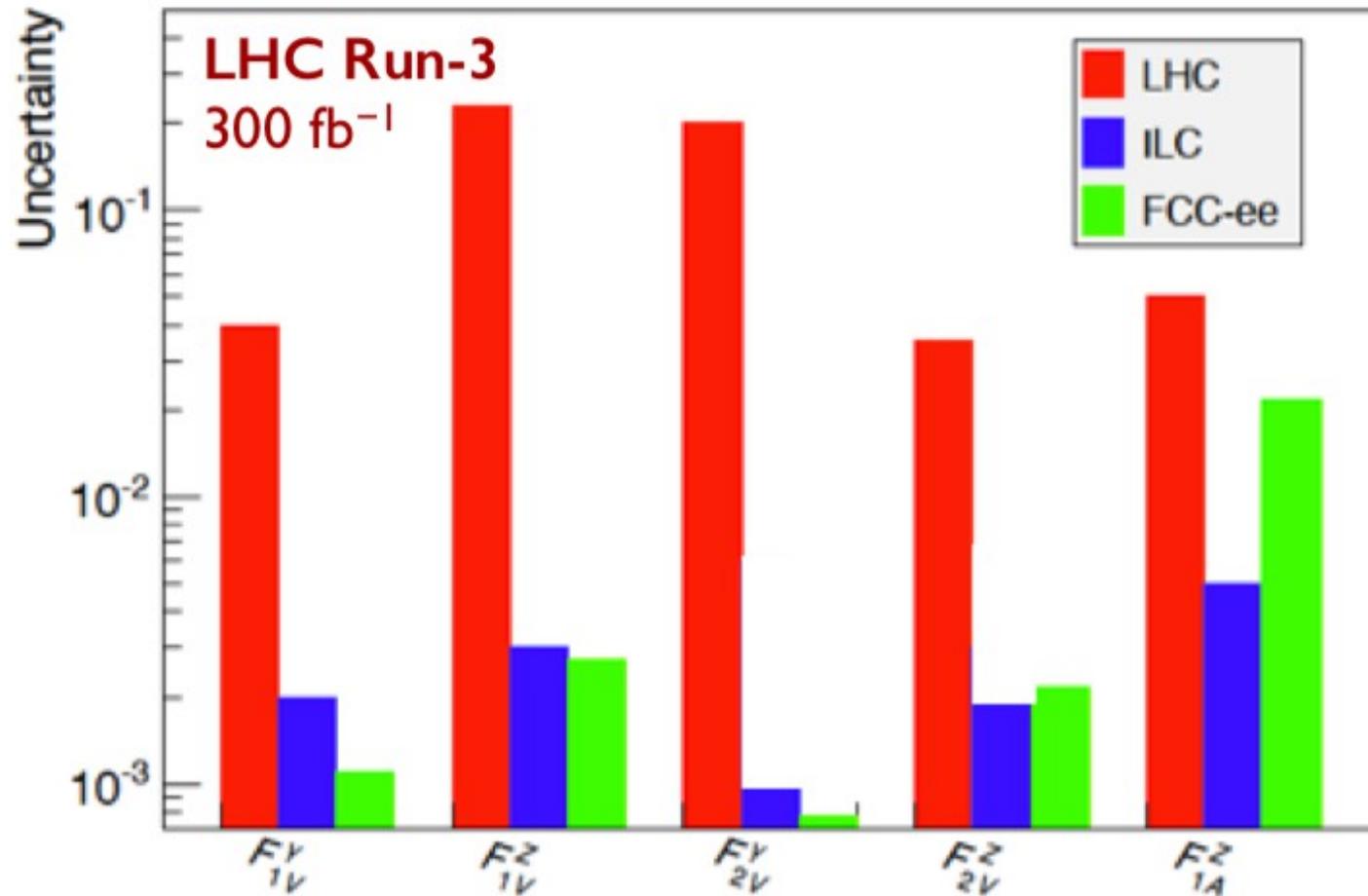
– Already by measuring 250 GeV

- Recall the sign uncertainty on LEP1 solutions
 $\delta g_{RZ}/g_{RZ} = 25\%$ or $\delta g_{RZ}/g_{RZ} = -225\%$

– Not a problem at 250 GeV to make the right choice for the sign due to interference photon/Z



Accuracy on CP conserving couplings



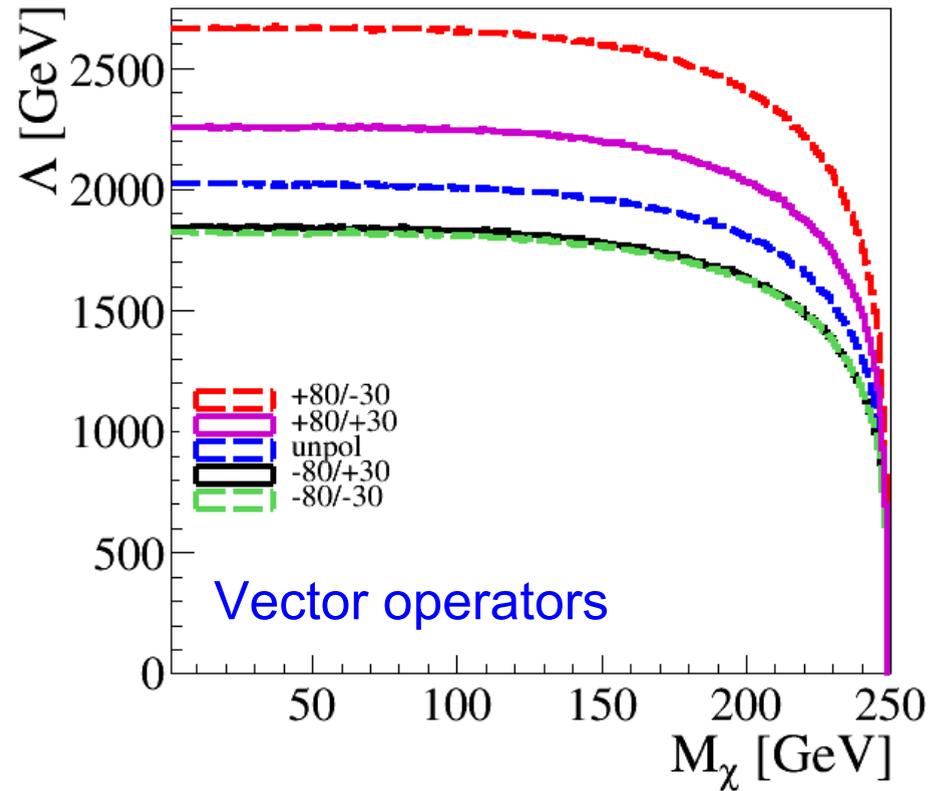
Arxiv:1503.01325

corrected for ILC values published in 1505.06020

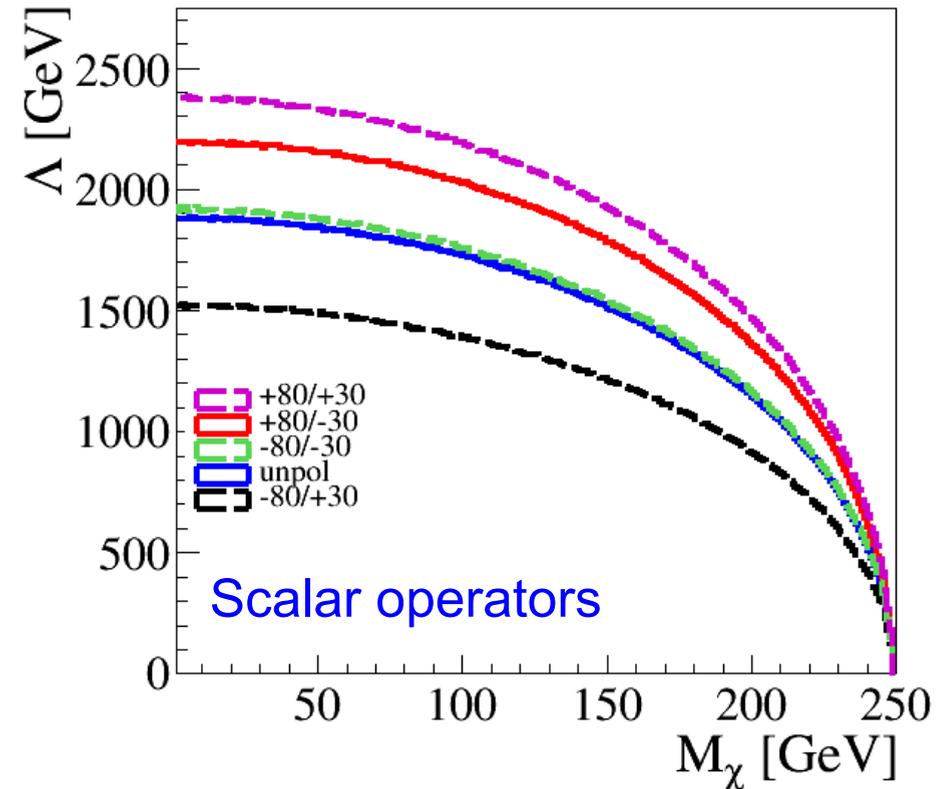
- e⁺e⁻ collider might be up to two orders of magnitude more precise than LHC ($\sqrt{s} = 14 \text{ TeV}$)
- Large disentangling of couplings for ILC thanks to polarised beams
- Final state analysis at FCCee
 - Also possible at LC => Redundancy
- Note
 - Maximal Lumi scenario for FCCee
 - Minimal Lumi scenario for ILC (~factor 4 possible with increased lumi and improved selection)

LC promises to be high precision machine for electroweak top couplings

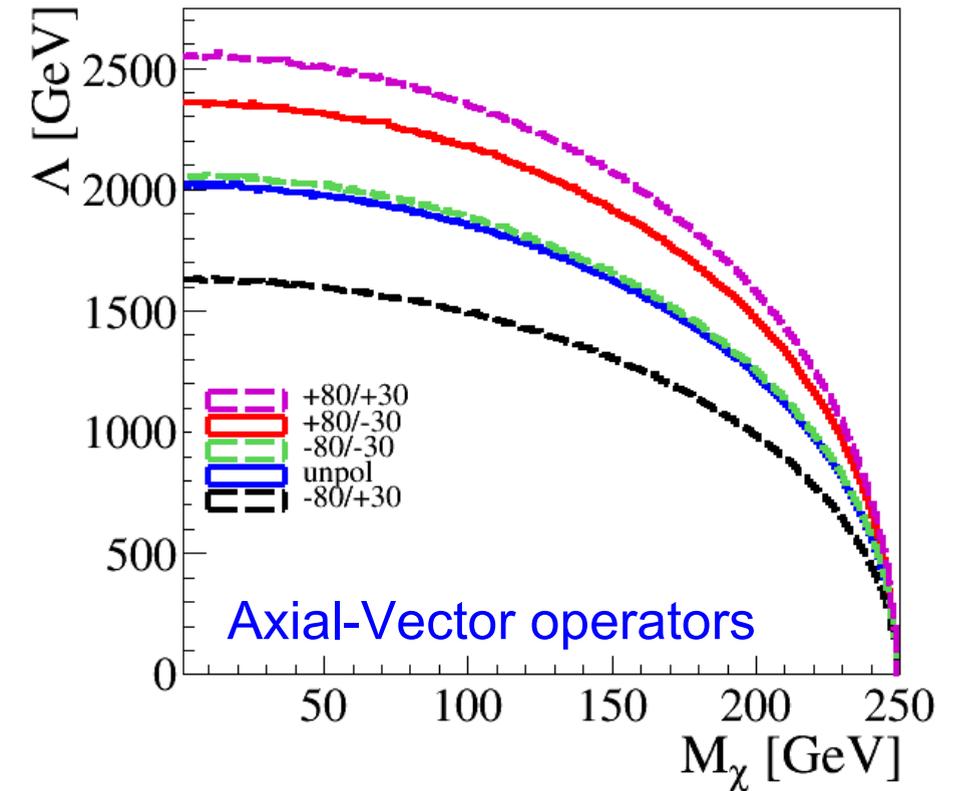
Vector operator, $\sqrt{s} = 500 \text{ GeV}$, 500 fb^{-1} , $3\sigma \text{ CL}$



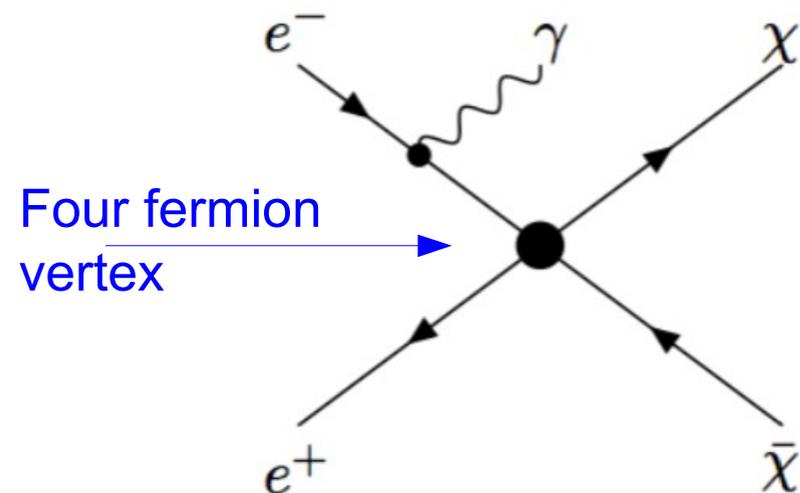
Scalar operator (s-channel), $\sqrt{s} = 500 \text{ GeV}$, 500 fb^{-1} , $3\sigma \text{ CL}$



Axial-vector operator, $\sqrt{s} = 500 \text{ GeV}$, 500 fb^{-1} , $3\sigma \text{ CL}$



M. Habermehl

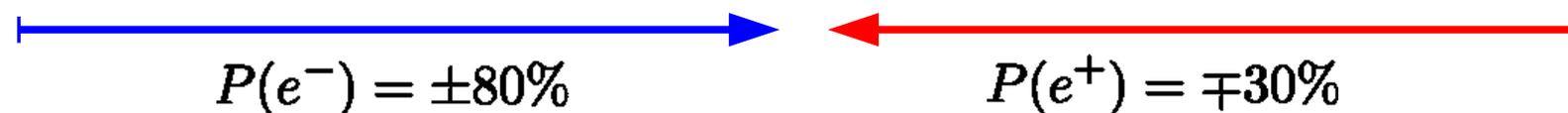


- Search (or set limits) for dark matter production
- Scale reach depends on beam polarisation
 - Largest scale reach for positive electron beam polarisation
 - Positron polarisation makes a difference

- **Beam polarisation is an essential asset for a successful e+e- precision program**
 - Remember that the SM is a chiral theory!!!
 - For comprehensive overviews see also hep-ph/0507011 and 1801.02840
- **Beam polarisation allows for large disentangling of various effects of new physics (or for constraining them further)**
 - Examples are electroweak fermion couplings and TGC
 - Helps a great deal to simplify analyses and interpretation of results due to adequate experimental setup for the theory under test
- **Linear Collider concept allows for sweeping over large energy for precision tests and **direct and indirect discoveries****
 - Measurements from Z pole to > 1 TeV within one facility
 - Colliders w/o strong beam polarisation will provide important complementary information
- **A clear pattern of anomalies would be an excellent (and maybe the only) motivation for a large hadron machine**

Backup

With two beam polarisation configurations



There exist a number of observables sensitive to chiral structure, e.g.

σ_I	$A_{FB,I}^t = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)}$	$(F_R)_I = \frac{(\sigma_{tR})_I}{\sigma_I}$
x-section	Forward backward asymmetry	Fraction of right handed top quarks



Extraction of relevant unknowns

$$\begin{aligned}
 &F_{1V}^\gamma, F_{1V}^Z, F_{1A}^\gamma = 0, F_{1A}^Z \quad \text{or equivalently} \quad g_L^\gamma, g_R^\gamma, g_L^Z, g_R^Z \\
 &F_{2V}^\gamma, F_{2V}^Z
 \end{aligned}$$

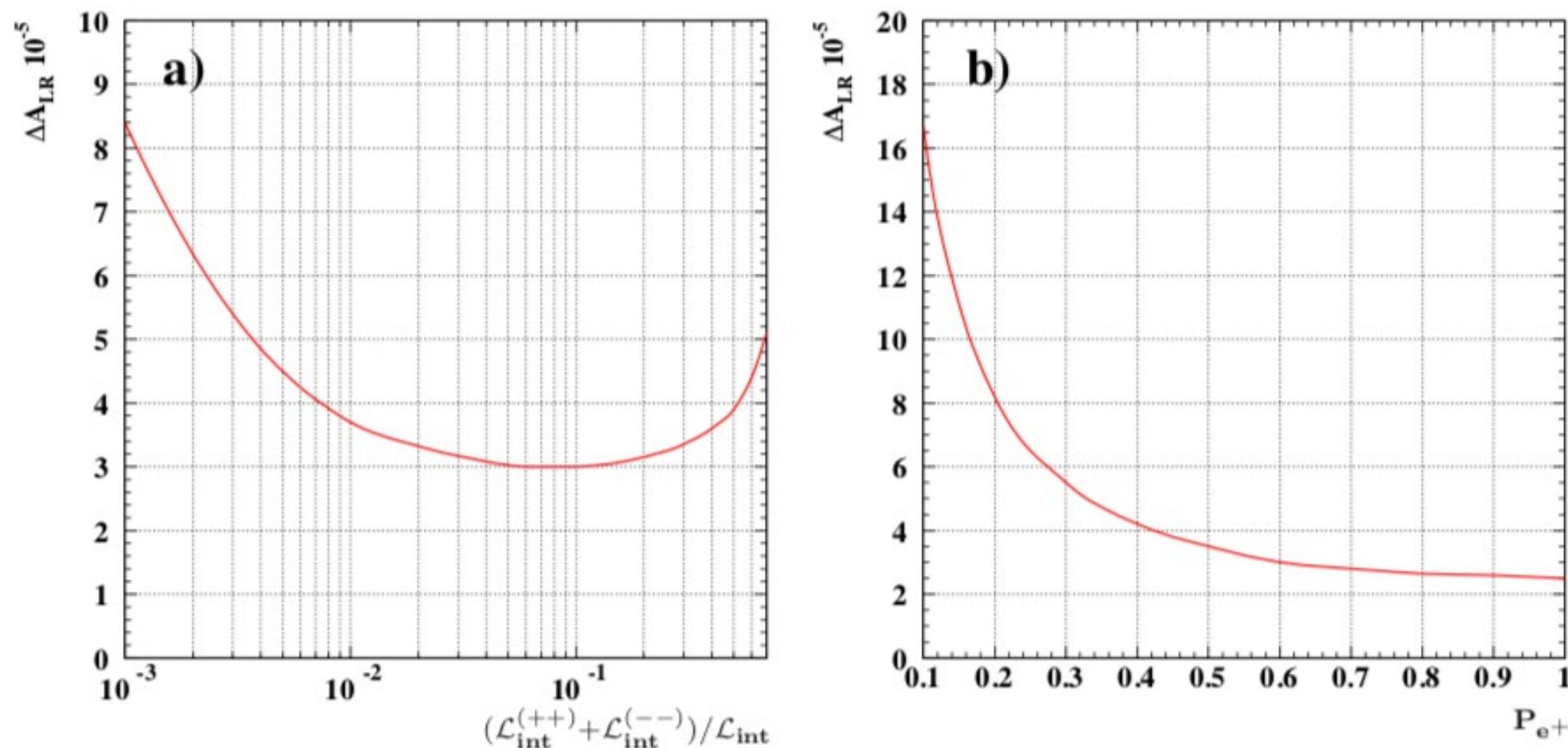


Figure 2.7: Test of the electroweak theory: the statistical error on A_{LR} of $e^+e^- \rightarrow Z \rightarrow \ell\bar{\ell}$ at GigaZ, (a) as a function of the fraction of luminosity spent on the less favoured polarization combinations σ_{++} and σ_{--} and (b) its dependence on P_{e+} for fixed $P_{e-} = \pm 80\%$ [51].

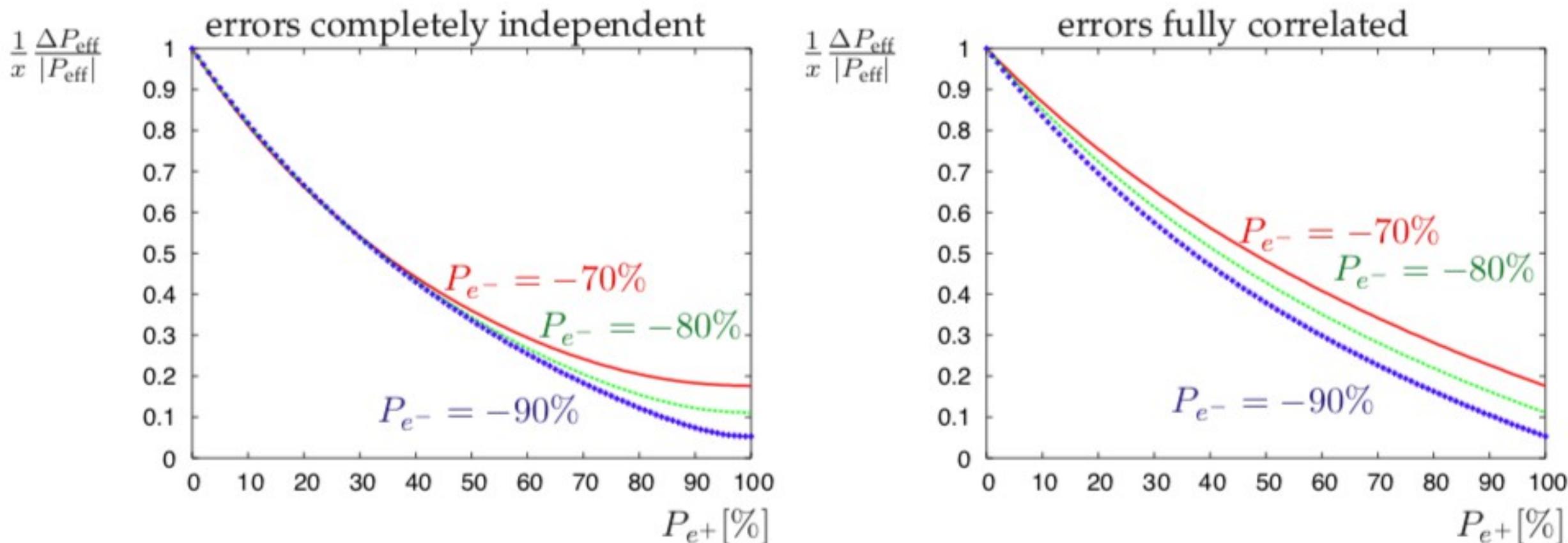


Figure 1.6: Relative uncertainty on the effective polarization, $\Delta P_{\text{eff}}/|P_{\text{eff}}| \sim \Delta A_{\text{LR}}/A_{\text{LR}}$, normalized to the relative polarimeter precision $x = \Delta P_{e^-}/P_{e^-} = \Delta P_{e^+}/P_{e^+}$ for independent and correlated errors on P_{e^-} and P_{e^+} , see eqs. (1.25), (1.27).

$$A_{\text{LR}} = \frac{1}{P_{\text{eff}}} A_{\text{LR}}^{\text{obs}} = \frac{1}{P_{\text{eff}}} \frac{\sigma_{-+} - \sigma_{+-}}{\sigma_{-+} + \sigma_{+-}}, \quad (1.24)$$



Helicity is projection of Spin $\vec{\sigma}$ onto direction $\hat{\mathbf{p}}$ of motion of massive particle

Eigenvalues of $\frac{1}{2}\vec{\sigma}\hat{\mathbf{p}}$:

-1/2  Left handed helicity

Caveat:
Helicity is frame dependent!

1/2  Right handed helicity

=> Not Lorentz invariant

Helicity projection operator

Chirality projection operator

$$\Pi^{\pm}(\mathbf{p}) = \frac{1}{2}(1 \pm \vec{\sigma}\hat{\mathbf{p}}) \xrightarrow{m=0} \Pi^{\pm}(\mathbf{p}) = \frac{1}{2}(1 \pm \gamma^5)$$

Chirality is projection of Spin $\vec{\sigma}$ onto direction $\hat{\mathbf{p}}$ of motion of massless particle

Chirality is frame independent! => Basis to define helicity states

$$u_L = \left(1 + \frac{|\vec{p}|}{E+m}\right) u_{LC} + \left(1 - \frac{|\vec{p}|}{E+m}\right) u_{RC} \quad E \gg m$$

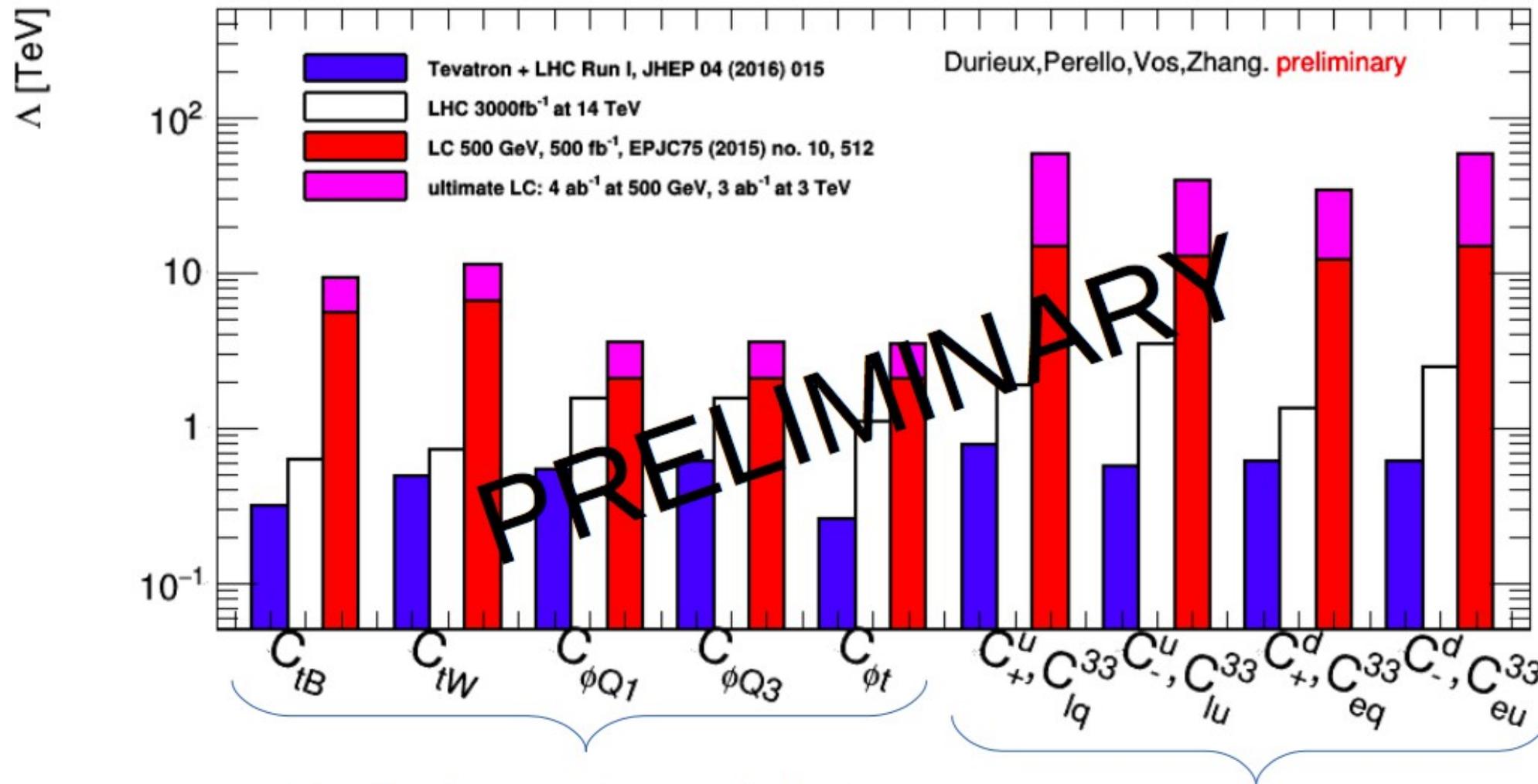
$$u_R = \left(1 - \frac{|\vec{p}|}{E+m}\right) u_{LC} + \left(1 + \frac{|\vec{p}|}{E+m}\right) u_{RC}$$

Tevatron + LHC from TopFitter (individual 95% limits)

Prospects for 3000/fb -> Schultz, Soreq, Vos, Perello ... + extrapolation

LC 500 prospects from arxiv: 1505.06020

Prospects somewhat speculative but may be covered by full ILC lumi



Two-fermion operators, equivalent to F1V, F1A, F2V form factors

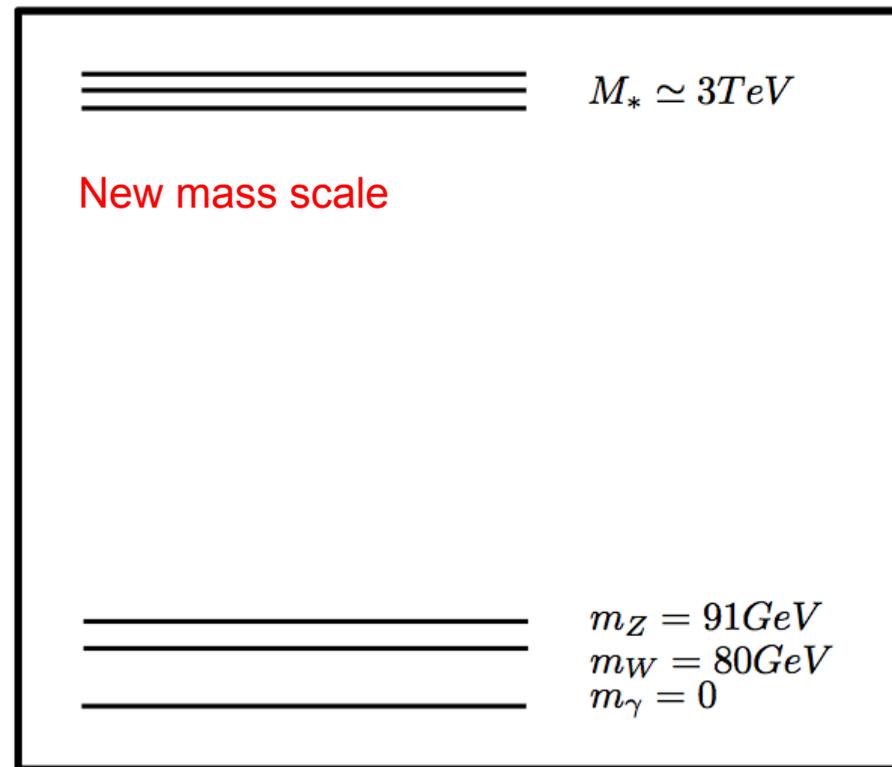
Four-fermion operators (qqtt at hadron colliders, eett at LC)

Compositeness:

- ... provides elegant solution for naturalness
- ... few tensions with SM predictions
- ... **all** scalar objects observed in nature turned out to be bound states of fermions
- ... Duality with Randall-Sundrum Models

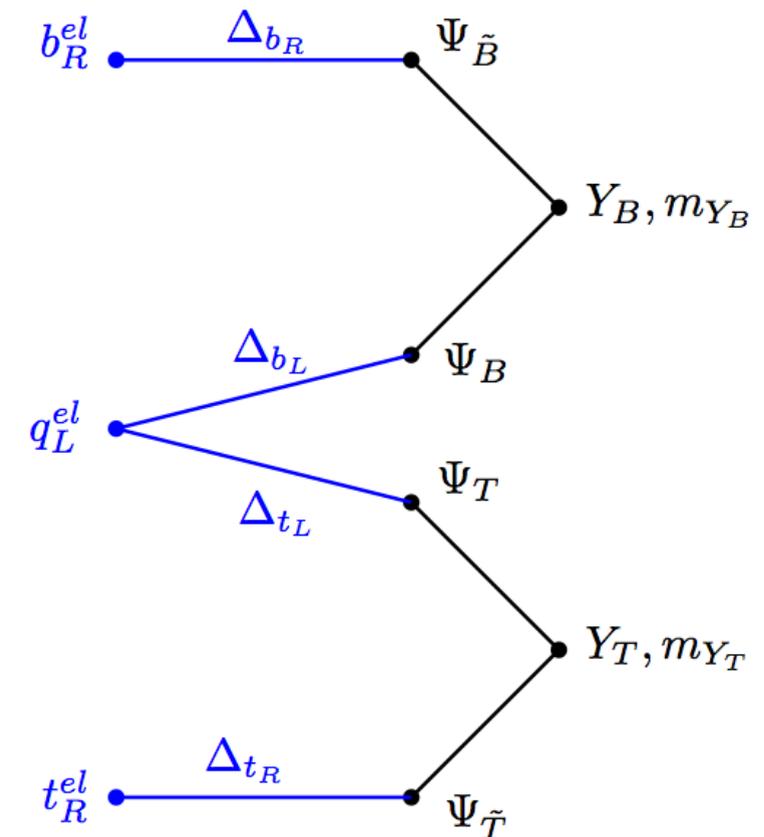
à la G.M. Pruna, LC 13, Trento

Bosonic sector mass spectrum



Fermionic resonances

From heavy left handed SM doublet and heavy right handed SM singlet



Physics modify Yukawa couplings and Z_{tt} , Z_{bb}
Heavy fermion effect!